

Chapter 5.10: Independent spins in a field: adiabatic demagnetization

We had for noninteracting Ising spins $H = -BS$ $Z = 2 \cosh \beta B$ $M = -\frac{\partial F}{\partial B} = \tanh \beta B$

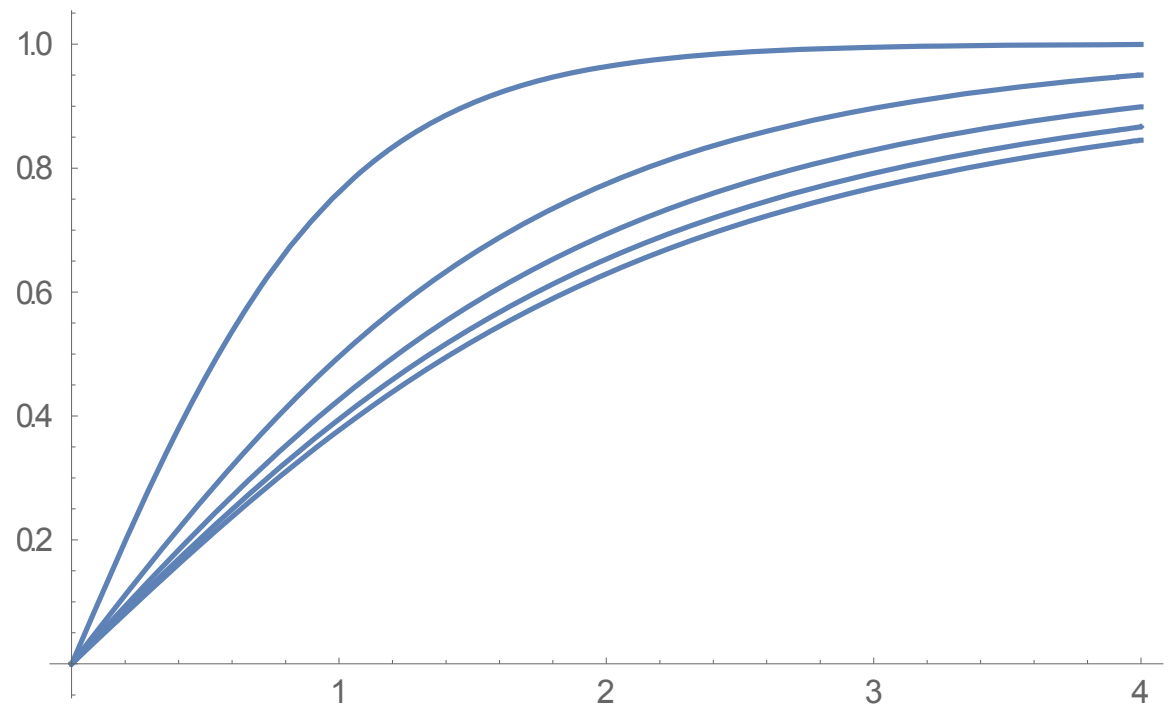
General paramagnetic coupling of field to magnetic moments:

$$H_{\text{field}} = -g\mu_B \vec{J} \cdot \vec{B} \quad \rightarrow \quad -Bj_z \quad \quad Z = \sum_{j_z=-j}^j e^{\beta B j_z} = \frac{\sinh \beta B (j + \frac{1}{2})}{\sinh \frac{\beta B}{2}}$$

Magnetization $M = -\frac{\partial F}{\partial B} = j B_j(j\beta B)$

Brillouin Function

$$B_j(x) = \frac{2j+1}{2j} \coth\left(\frac{2j+1}{2j}x\right) - \frac{1}{2j} \coth\left(\frac{x}{2j}\right)$$



5.10-2 Independent Spins in a field: adiabatic demagnetization

Susceptibility: Curie law

General: $\chi = \frac{\partial M}{\partial B} = -\frac{\partial^2 F}{\partial B^2}$

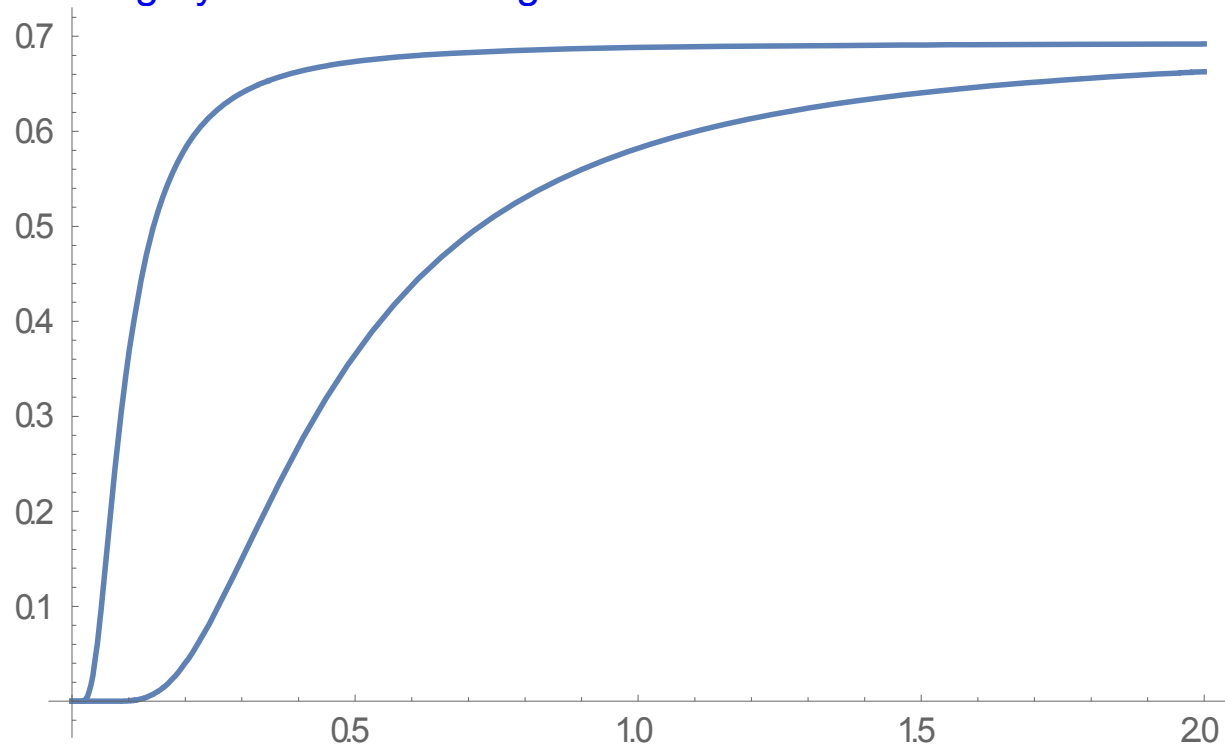
Small fields: $B_j(x) \rightarrow \frac{j+1}{3j} x$

5.10-3 Independent Spins in a field: adiabatic demagnetization

Entropy

$$S = -\frac{\partial F}{\partial T} = k_B \ln Z + \frac{k_B T}{Z} \frac{\partial Z}{\partial T}$$

Cooling by adiabatic demagnetization



Third law of thermodynamics

Nernst 1912: "It is impossible for any procedure to lead to the isotherm $T = 0$ in a finite number of steps."

