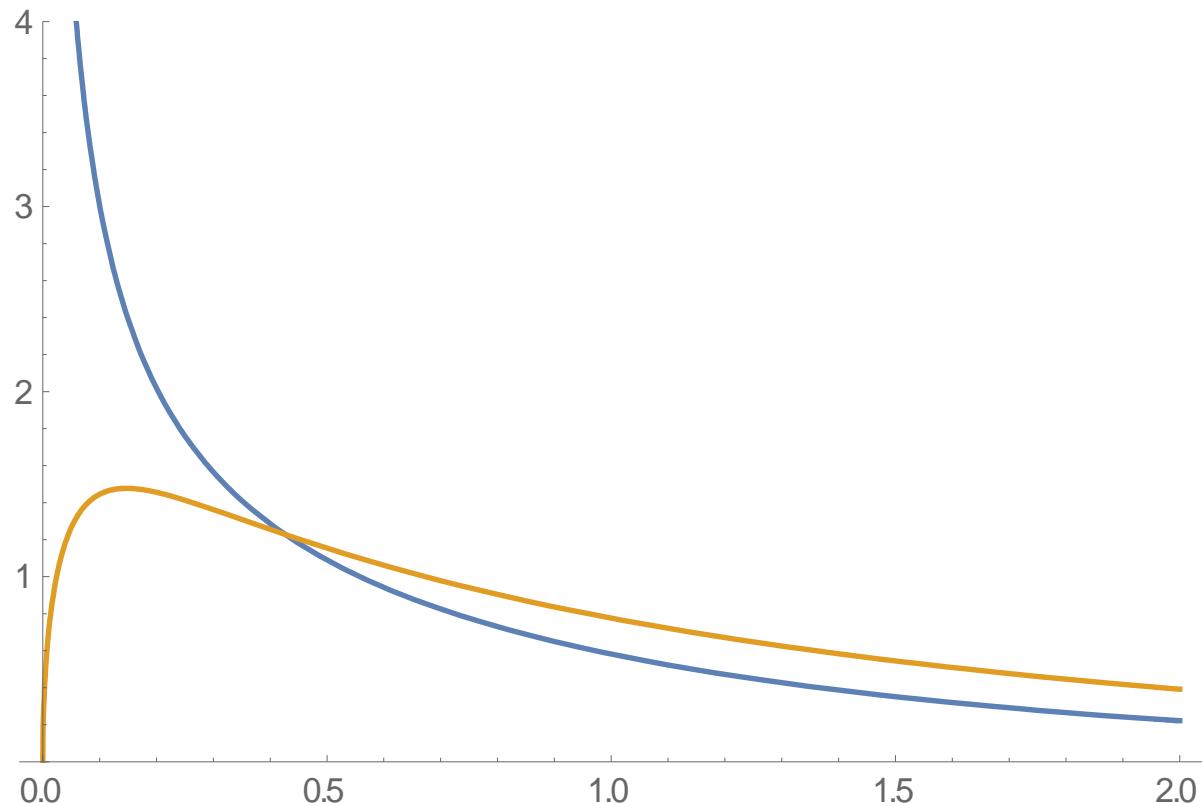


## Chapter 4.7: The Bose Einstein Condensation

$$N = \sum_r \langle n_r \rangle \rightarrow \int_0^\infty d\varepsilon \frac{g(\varepsilon)}{z^{-1} e^{\beta\varepsilon} - 1} = \frac{V}{\lambda_T^3} \text{Li}_{3/2}(z)$$

$$\varepsilon = (n_x^2 + n_y^2 + n_z^2) \frac{\hbar^2 \pi^2}{2mL^2}$$



#### 4.7-2 The Bose Einstein Condensation

Corrected number calculation: ground state separately

$$\lambda_T = \hbar \sqrt{\frac{2\pi}{mk_B T}}$$

$$N = \sum_r \langle n_r \rangle \rightarrow \langle n_0 \rangle + \int_0^\infty d\varepsilon \frac{g(\varepsilon)}{z^{-1} e^{\beta\varepsilon} - 1} = \frac{1}{z^{-1} - 1} + \frac{V}{\lambda_T^3} \text{Li}_{3/2}(z)$$

$$x_0 = \frac{\langle n_0 \rangle}{N} = \frac{1}{N} \frac{1}{z^{-1} - 1}$$

becomes finite as  $z \rightarrow 1$  and  $N \rightarrow \infty, V \rightarrow \infty$

Temperature at which ground state occupation  $x_0$  becomes significant is

$$N = \frac{V}{\lambda_{T_c}^3} \text{Li}_{3/2}(1) = \frac{V}{\lambda_{T_c}^3} \zeta\left(\frac{3}{2}\right)$$

#### 4.7-3 The Bose Einstein Condensation

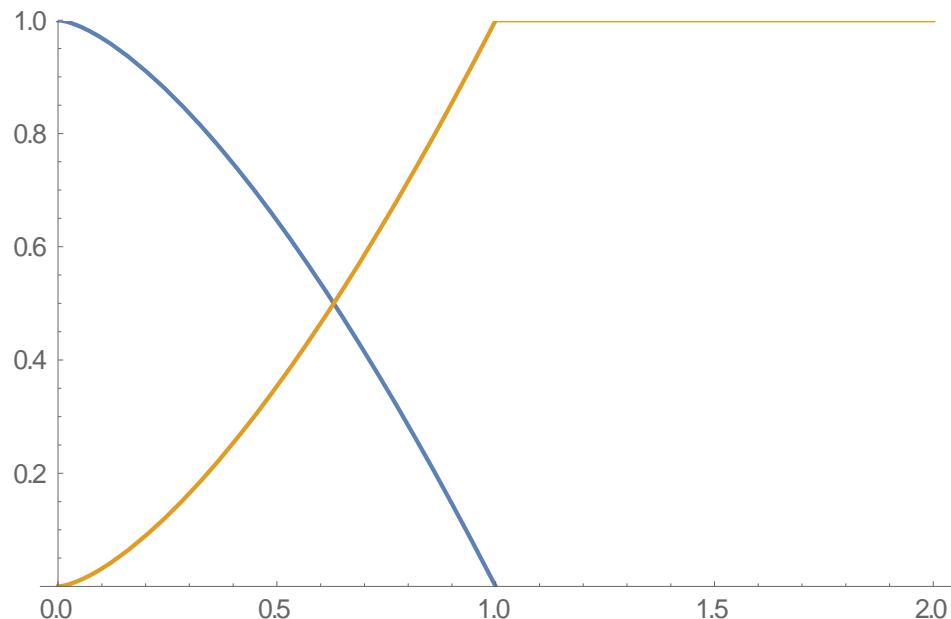
For  $T > T_c$ :

$$N = \frac{V}{\lambda_T^3} \text{Li}_{3/2}(z)$$

$$\frac{V}{N} = \frac{\lambda_{T_c}^3}{\zeta(3/2)}$$

For  $T < T_c$ :

$$N = \frac{1}{z^{-1} - 1} + \frac{V}{\lambda_T^3} \text{Li}_{3/2}(1)$$



#### 4.7-4 The Bose Einstein Condensation

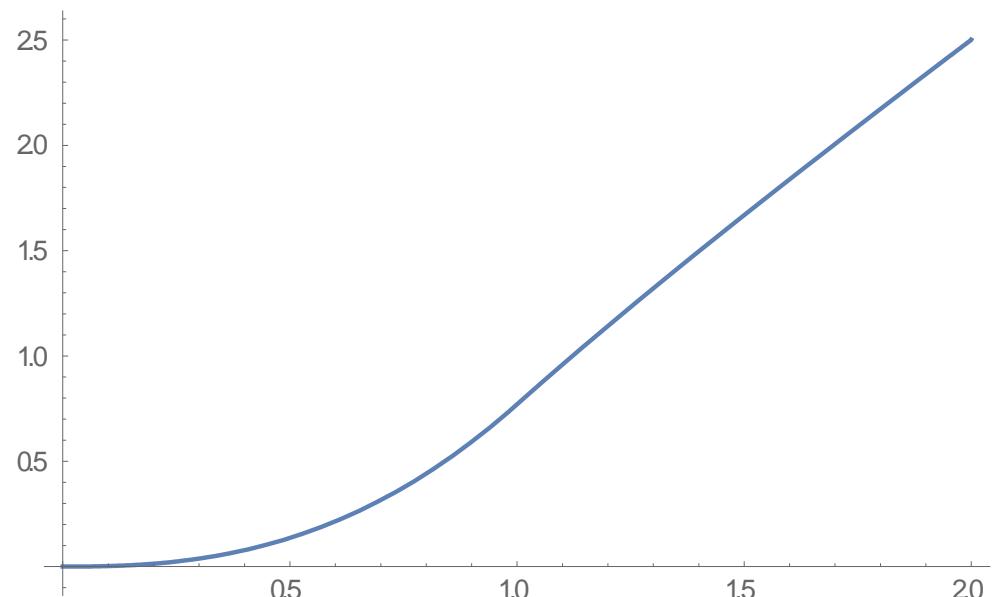
Consequences for energy and pressure

$$\frac{V}{N} = \frac{\lambda_{T_c}^3}{\zeta(3/2)}$$

$$E = \sum_r \varepsilon_r \langle n_r \rangle = \frac{\varepsilon_0}{z^{-1} - 1} + \int_0^\infty d\varepsilon \frac{\varepsilon g(\varepsilon)}{z^{-1} e^{\beta\varepsilon} - 1} = \frac{3}{2} \frac{k_B T V}{\lambda_T^3} \text{Li}_{5/2}(z)$$

$$p = - \left( \frac{\partial \Phi}{\partial V} \right)_{\alpha, T} \rightarrow \frac{k_B T}{\lambda_T^3} \text{Li}_{5/2}(z)$$

#### 4.7-5 The Bose Einstein Condensation



#### 4.7-6 The Bose Einstein Condensation

##### Consequences for specific heat

$$c_V = - \left( \frac{\partial E}{\partial T} \right)_V$$

