

## Chapter 3.3: The Ideal Gas in the Canonical Ensemble

Independent particles  $\{\vec{r}_1, \vec{p}_1, \vec{r}_2, \vec{p}_2, \vec{r}_3, \vec{p}_3, \dots, \vec{r}_N, \vec{p}_N\}$

$$\int_{-\infty}^{\infty} dx e^{-\alpha x^2} = \sqrt{\frac{\pi}{\alpha}}$$

$$\varepsilon(\{\vec{r}_j, \vec{p}_j\}) = \sum_j \frac{\vec{p}_j^2}{2m} = \sum_j \varepsilon_j$$

$$\varepsilon_j = \frac{p_j^2}{2m}$$

$$\begin{aligned} Z &= (Z_1)^N \\ &= V^N \left(\frac{2\pi m}{\beta}\right)^{\frac{3}{2}N} \end{aligned}$$

$$Z_1 = \int d\vec{r}_1 d\vec{p}_1 e^{-\beta \varepsilon_1} = V \int d\vec{p} e^{-\beta \frac{p^2}{2m}}$$

$$= V \int dp_x e^{-\beta \frac{p_x^2}{2m}} \int dp_y e^{-\beta \frac{p_y^2}{2m}} \int dp_z e^{-\beta \frac{p_z^2}{2m}}$$

$$\ln Z = N \left( \ln V - \frac{3}{2} \ln \beta + \text{const} \right)$$

$$= V \left(\frac{2\pi m}{\beta}\right)^{\frac{3}{2}}$$

$$E = -\frac{\partial \ln Z}{\partial \beta} = +N \frac{3}{2} \frac{1}{\beta} = \frac{3}{2} N k_B T \quad \checkmark$$

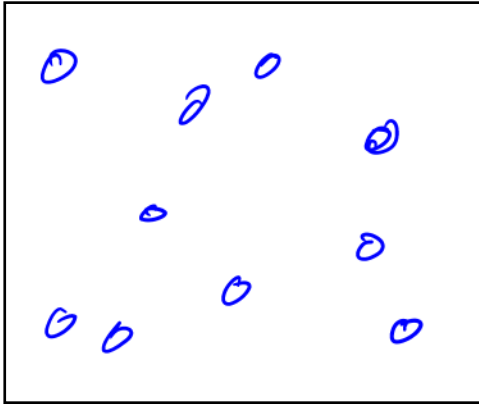
$$\frac{E}{T} = \frac{3}{2} N k_B$$

$$p = \left(\frac{\partial F}{\partial V}\right)_T = k_B T \left(\frac{\partial \ln Z}{\partial V}\right)_T = \frac{N k_B T}{V} \quad \checkmark$$

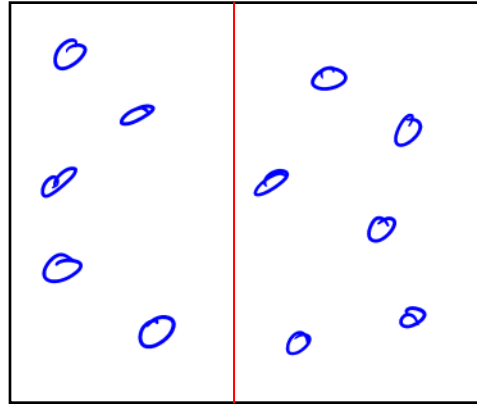
$$* \quad S = -\frac{F}{T} + \frac{E}{T} = k_B \ln Z + \frac{E}{T} = N k_B \left( \ln V - \frac{3}{2} \ln \beta + \text{const} \right)$$

3.3-2 The Ideal Gas in the Canonical Ensemble

Entropy



$N/2, V/2$        $N/2, V/2$



$S_{V/2}$        $S_{V/2}$

$$S_{\text{before}} = N k_B \left( \ln \frac{V}{N} - \frac{3}{2} \ln \beta + \text{const} \right)$$

$$S_{\text{after}} = S_{V/2} + S_{V/2} = 2 \times S_{V/2}$$

$$= 2 \times \frac{N}{2} k_B \left( \ln \frac{V/2}{N/2} - \frac{3}{2} \ln \beta + \text{const} \right)$$

$$= N k_B \left( \ln V - \frac{3}{2} \ln \beta + \text{const} \right) - N k_B \ln 2$$

↙ reversible!

$$\underline{S_{\text{after}}} = S_{\text{before}} - N k_B \ln 2$$

3.3-3 The Ideal Gas in the Canonical Ensemble

Gibbs paradox and Gibbs factor

$\frac{1}{N!}$

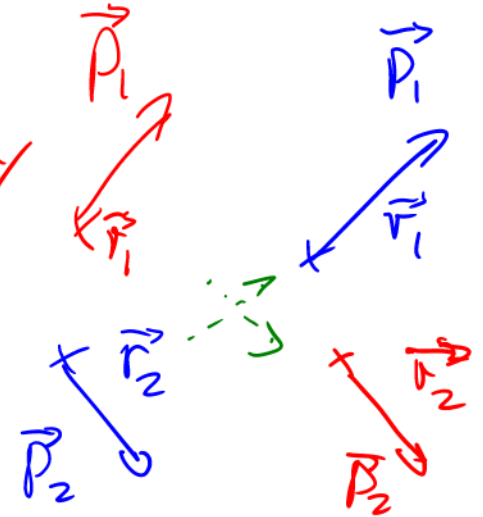
Same state

For indistinguishable particles avoid overcounting by exchanging indices:

integration over all  $\{\vec{p}_j, \vec{r}_j\}$  there are too many microstate

$$\int \prod_j d^3\vec{r}_j d^3\vec{p}_j \rightarrow \frac{1}{N!} \int \prod_j d^3\vec{r}_j d^3\vec{p}_j$$

$$Z = \frac{(z_1)^N}{N!} = \frac{1}{N!} \int d\vec{r} e^{-\beta \sum_j \epsilon_j}$$

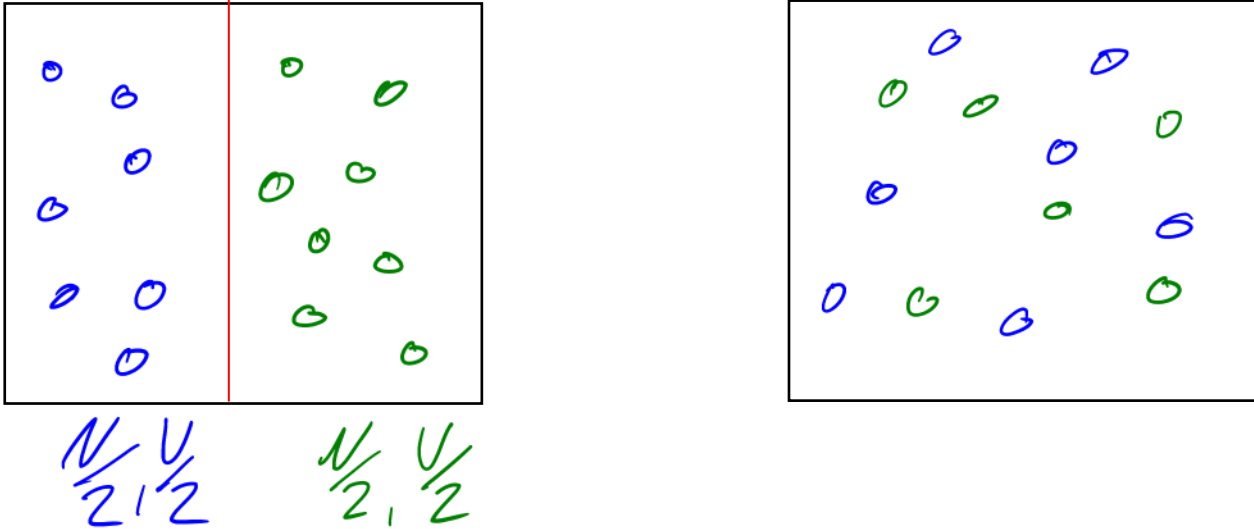


$$\ln Z = N \ln z_1 - \ln N! = N \ln z_1 - N \ln N + N$$

$$S = k_B \ln Z + \frac{E}{T} = k_B N \left( \ln V - \frac{3}{2} \ln \beta + \text{const} \right) - k_B N \ln N$$

$$= k_B N \left( \ln \frac{V}{N} - \frac{3}{2} \ln \beta + \text{const} \right)$$

## Entropy of mixing



$$S_{\text{before}} = 2 \times \frac{N}{2} k_B \left( \ln \frac{V/2}{2 \cdot N/2} - \frac{3}{2} \ln \beta + \text{const} \right)$$

$$S_{\text{after}} = 2 \times \frac{N}{2} k_B \left( \ln \frac{2V}{N} - \frac{3}{2} \ln \beta + \text{const} \right) = S_{\text{before}} + k_B N \ln 2$$

"Entropy of mixing" non-reversible