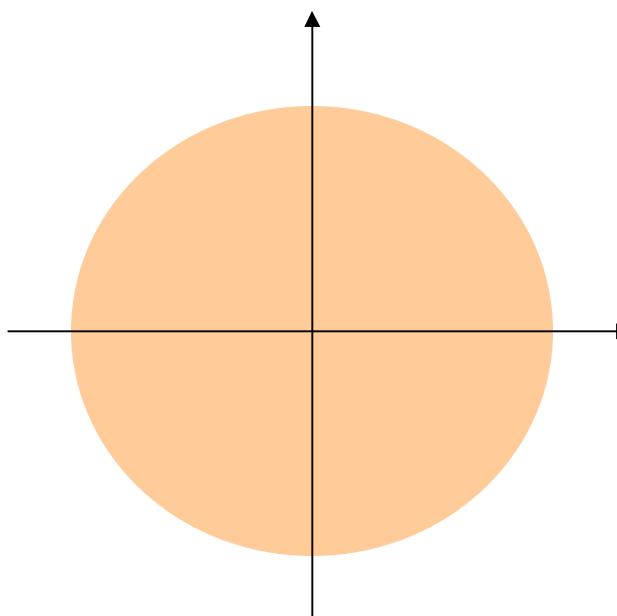


Stability of the excitations

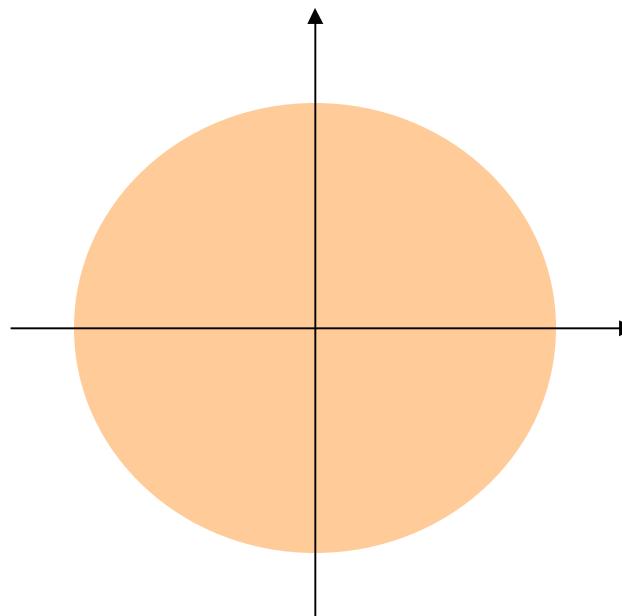
Decay probability using two constraints

Inverse lifetime is proportional



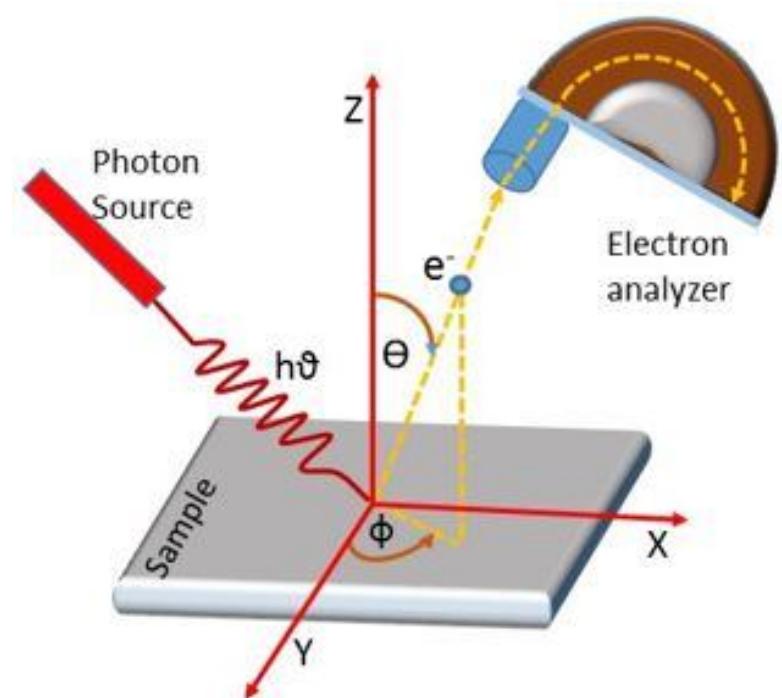
Effective lifetime and correlations

$$c_{\sigma}^{\dagger}(\vec{k}, t) = e^{-it\varepsilon_{\vec{k}}/\hbar} c_{\sigma}^{\dagger}(\vec{k})$$



Spectral density: ARPES

$$\rho(\omega, \vec{k}) = \sum_{\lambda} \left| \langle \lambda | c_{\sigma}(\vec{k}) | GS \rangle \right|^2 \delta(\omega - \omega_{\lambda})$$



Retarded Green's Function

$$\rho(\omega, \vec{k}) = -\frac{1}{\pi} \text{Im} G_R(\omega, \vec{k}) \quad \text{where} \quad G_R(\omega, \vec{k}) = -i \int_0^\infty dt e^{i\omega t} \langle \{c_\sigma(\vec{k}, t), c_\sigma^\dagger(\vec{k}, 0)\} \rangle$$

Self Energy and quasi-particle spectral weight

$$c_\sigma(\vec{k}, t) = e^{i\omega_{\vec{k}} t} e^{-i\Gamma_{\vec{k}} t} c_\sigma(\vec{k})$$

$$\rho(\omega, \vec{k}) = -\frac{1}{\pi} \text{Im} G_R(\omega, \vec{k})$$

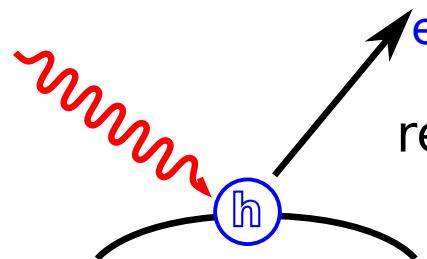
$$G_R(\omega, \vec{k}) = \frac{Z}{\omega - \omega_{\vec{k}} - \Sigma_{\vec{k}}}$$

Spectral functions

Spectral function

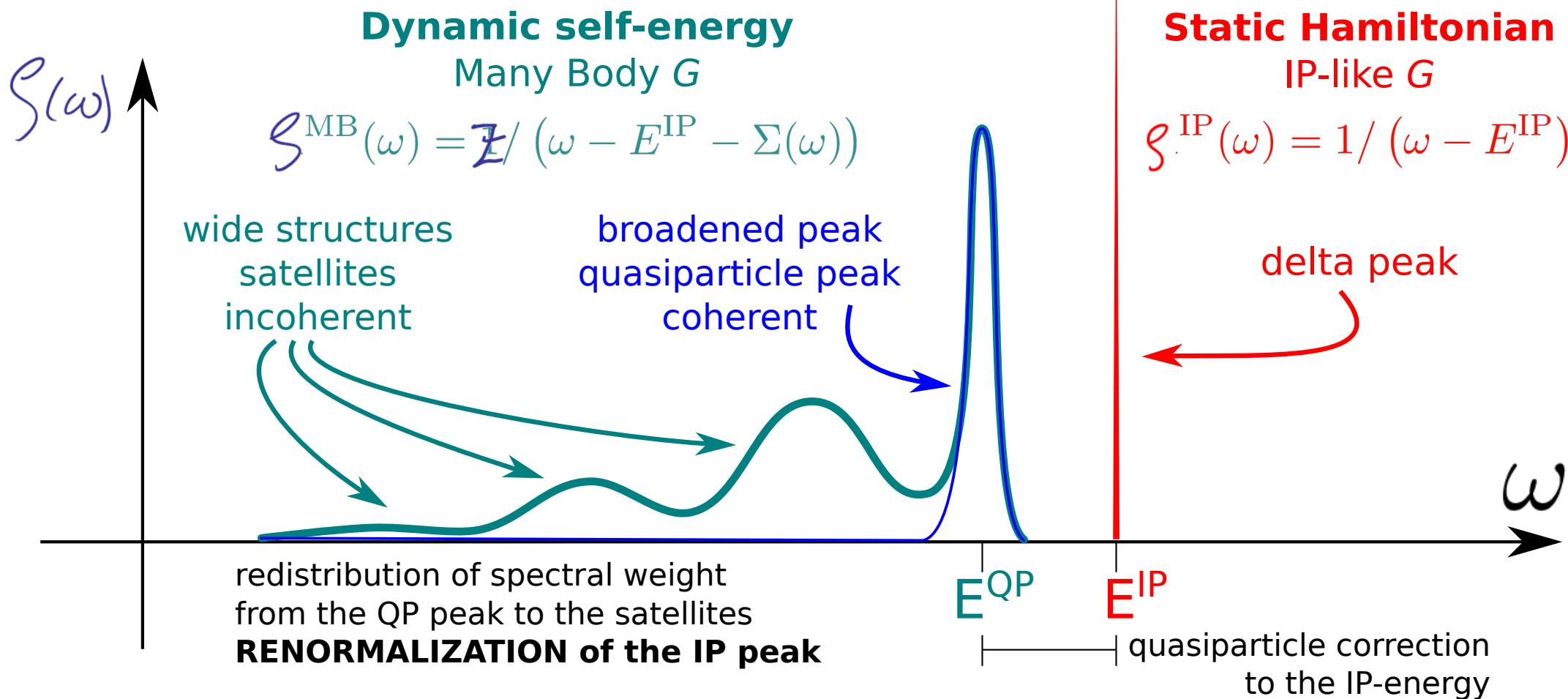
structures at poles of G

$$\mathcal{S}(\omega) = \frac{1}{\pi} |\text{Im}G(\omega)|$$



Photoemission

removal energy (hole creation)



General Mean Field Theory

$$\text{e.g. } \hat{H}_{\text{int}} = \sum_{\sigma_1, \sigma_2} \int d^3 \vec{r}_1 d^3 \vec{r}_2 \hat{\Psi}_{\sigma_1}^\dagger(\vec{r}_1) \hat{\Psi}_{\sigma_2}^\dagger(\vec{r}_2) V(\vec{r}_1 - \vec{r}_2) \hat{\Psi}_{\sigma_2}(\vec{r}_2) \hat{\Psi}_{\sigma_1}(\vec{r}_1)$$

or Hubbard interaction $U n_{j,\uparrow} n_{j,\downarrow}$

$$\hat{H}_{\text{int}} = \hat{A} \hat{B}$$

Hartree Fock Theory

Variational ansatz for a product state with N particles $|\psi_{HF}\rangle = \Psi_1^\dagger \Psi_2^\dagger \Psi_3^\dagger \dots \Psi_N^\dagger |0\rangle$

where an arbitrary single particle wavefunction is described by $\Psi_1^\dagger |0\rangle = \int d^3\vec{r} \varphi_1(\vec{r}) \hat{\Psi}^\dagger(\vec{r}) |0\rangle$

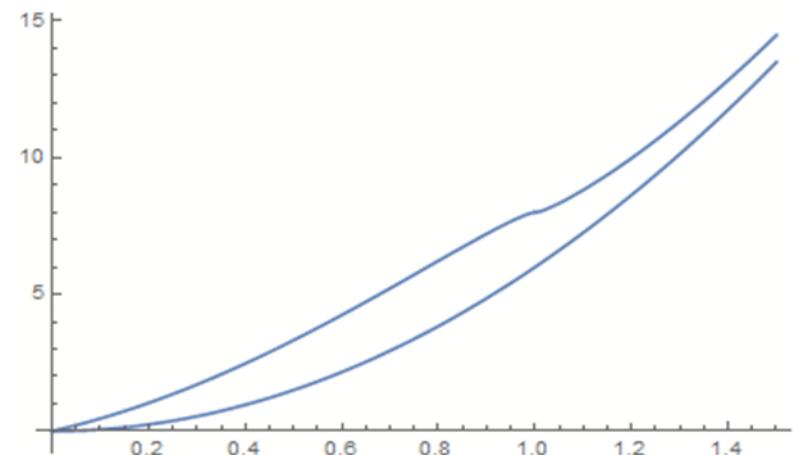
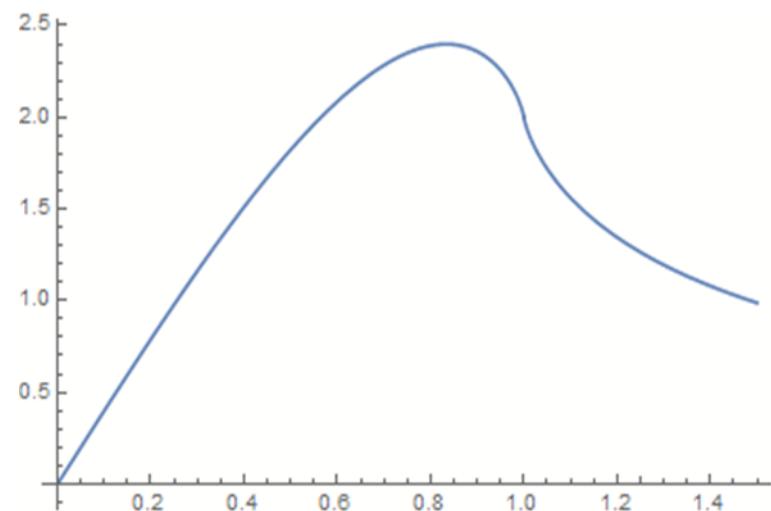
General two-body interaction $\hat{H}_{\text{int}} = \sum \hat{\Psi}_a^\dagger \hat{\Psi}_b^\dagger V \hat{\Psi}_c \hat{\Psi}_d$

Hartree and Fock terms as a mean fields:

$$c_{\sigma_1}^\dagger(\vec{k}_1)c_{\sigma_2}^\dagger(\vec{k}_2)\tilde{V}(\Delta\vec{k})c_{\sigma_2}(\vec{k}_2 - \Delta\vec{k})c_{\sigma_1}(\vec{k}_1 + \Delta\vec{k})$$

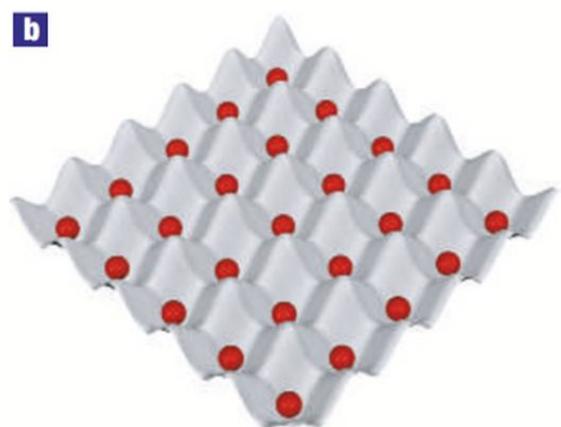
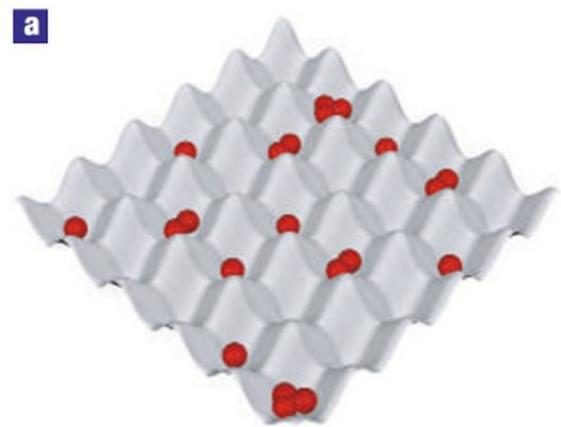
Example: Coulomb gas

$$\tilde{V}(\Delta\vec{k}) = \int d^3\vec{r} e^{i\vec{r} \cdot \Delta\vec{k}} \frac{e^2}{(2\pi)^3 r} = \frac{e^2}{2\pi^2 |\Delta\vec{k}|^2}$$



The Bose Hubbard model

$$\hat{H}_{BH} = -t \sum_{\langle i,j \rangle} \hat{a}_i^\dagger \hat{a}_j + U \sum_j \hat{n}_j (\hat{n}_j - 1) - \mu \sum_j \hat{n}_j$$



Off Diagonal Long Range Order

$$G_1(i - j) = \langle \hat{a}_i^\dagger \hat{a}_j \rangle$$

Mean Field Ansatz

$$\hat{a}_i^\dagger \hat{a}_j = \hat{a}_i^\dagger \langle \hat{a}_j \rangle + \langle \hat{a}_i^\dagger \rangle \hat{a}_j - \langle \hat{a}_i^\dagger \rangle \langle \hat{a}_j \rangle$$

$$\langle \hat{a}_j \rangle = \psi$$

$$\hat{H}_{BH} = -t \sum_{\langle i,j \rangle} \hat{a}_i^\dagger \hat{a}_j + U \sum_j \hat{n}_j (\hat{n}_j - 1) - \mu \sum_j \hat{n}_j$$

Mean Field Calculation

$$\hat{H}_{MF} = -t \sum_{\langle i,j \rangle} \left(\hat{a}_i^\dagger \psi + \psi^* \hat{a}_j - |\psi|^2 \right) + U \sum_j \hat{n}_j (\hat{n}_j - 1) - \mu \sum_j \hat{n}_j$$

Condition at phase transition: $\langle \hat{a}_j \rangle = \psi \rightarrow 0$

$$\hat{H}_0 = U\hat{n}(\hat{n} - 1) - \mu\hat{n}$$

$$\hat{H}_1 = -tz(\hat{a}^\dagger\psi + \psi^*\hat{a} - |\psi|^2)$$

Phase transition line

$$zt = \frac{(\mu - nU)((n-1)U - \mu)}{U + \mu}$$

Boson localization and the superfluid-insulator transition

Matthew P. A. Fisher, Peter B. Weichman, G. Grinstein, and Daniel S. Fisher
Phys. Rev. B 40, 546 – Published 1 July 1989