

## Summary of electron in a periodic potential: different basis systems

$$\hat{H} = \int d^3\vec{r} \hat{\Psi}^\dagger(\vec{r}) \left( -\frac{\hbar^2}{2m} \nabla^2 + U(\vec{r}) \right) \hat{\Psi}(\vec{r})$$

$$\begin{aligned} \hat{H} &= \hat{H}_{kin} + \hat{H}_{pot} = \frac{\hbar^2}{2m} \int d^3\vec{k} |\vec{k}|^2 \hat{c}^\dagger(\vec{k}) \hat{c}(\vec{k}) + \int d^3\vec{k} \sum_{\vec{G}} \hat{c}^\dagger(\vec{k} + \vec{G}) U(\vec{G}) \hat{c}(\vec{k}) \\ &= \sum_{\vec{G}} \int_{1BZ} d^3\vec{k} \left( \frac{\hbar^2}{2m} |\vec{k} - \vec{G}|^2 \hat{c}^\dagger(\vec{k} - \vec{G}) \hat{c}(\vec{k} - \vec{G}) + \sum_{\vec{G}'} \hat{c}^\dagger(\vec{k} - \vec{G} + \vec{G}') U(\vec{G}') \hat{c}(\vec{k} - \vec{G}) \right) \end{aligned}$$

$$\hat{H}_{\vec{k}} = \sum_{\alpha} \varepsilon_{\alpha}(\vec{k}) \hat{\Psi}_{\vec{k},\alpha}^\dagger \hat{\Psi}_{\vec{k},\alpha}$$

$$\hat{H} = \int_{1BZ} d^3\vec{k} \hat{H}_{\vec{k}} = \sum_{\alpha} \sum_{\vec{R}, \Delta\vec{R}} t_{\alpha}(\Delta\vec{R}) \tilde{\Psi}_{\vec{R}+\Delta\vec{R},\alpha}^\dagger \tilde{\Psi}_{\vec{R},\alpha}$$

$$\hat{c}^\dagger(\vec{k}) = \frac{1}{(2\pi)^{3/2}} \int d^3\vec{r} e^{i\vec{k}\cdot\vec{r}} \hat{\Psi}^\dagger(\vec{r})$$

$$U(\vec{G}) = \frac{1}{V_U} \int_{V_U} d^3\vec{r} e^{-i\vec{G}\cdot\vec{r}} U(\vec{r})$$

$$\Psi_{\vec{k},\alpha}^\dagger = \sum_{\vec{G}} \lambda_{\vec{k},\alpha}(\vec{G}) \hat{c}^\dagger(\vec{k} - \vec{G})$$

$$\tilde{\Psi}_{\vec{R},\alpha}(\vec{r}) = \frac{1}{\sqrt{V_{1BZ}}} \int_{V_{1BZ}} d^3\vec{k} e^{-i\vec{k}\cdot\vec{R}} \Psi_{\vec{k},\alpha}^\dagger$$

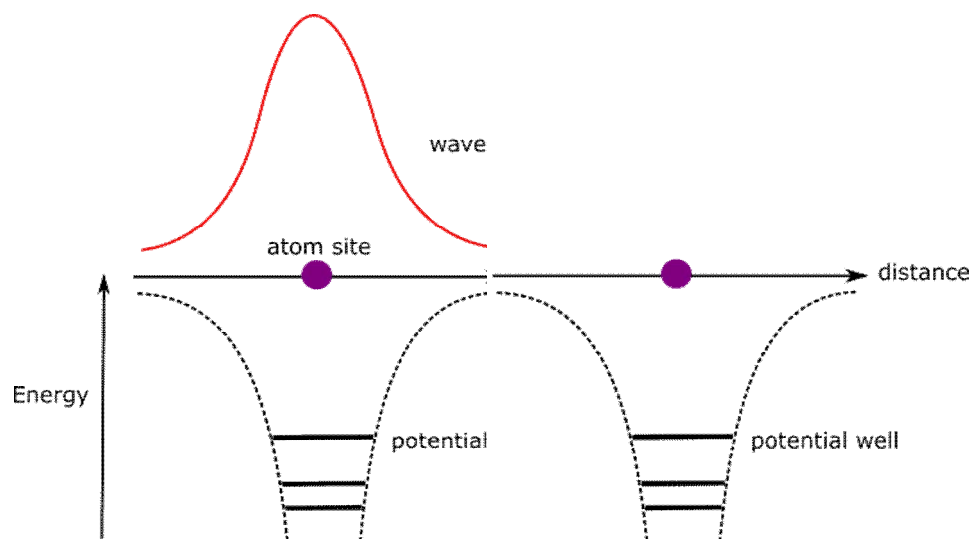
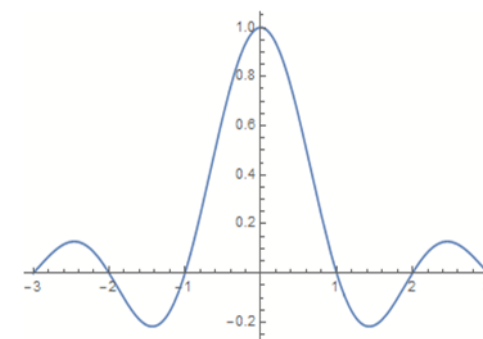
$$t_{\alpha}(\Delta\vec{R}) = \frac{1}{V_{1BZ}} \int_{V_{1BZ}} d^3\vec{k} \varepsilon_{\alpha}(\vec{k}) e^{i\vec{k}\cdot\Delta\vec{R}}$$

## Tight binding approximation

Hopping integrals  $t_\alpha(\Delta\vec{R})$  are Fourier transformation of dispersion but also matrix elements:

$$t_\alpha(\Delta\vec{R}) = \left\langle \tilde{\psi}_{\vec{R}+\Delta\vec{R},\alpha} \left| \hat{H} \right| \tilde{\psi}_{\vec{R},\alpha} \right\rangle$$

$$\tilde{\psi}_{\vec{R},\alpha}(\vec{r}) = \frac{1}{\sqrt{V_{1BZ}}} \int_{V_{1BZ}} d^3\vec{k} e^{-i\vec{k}\cdot\vec{R}} \psi_{\vec{k},\alpha}(\vec{r})$$

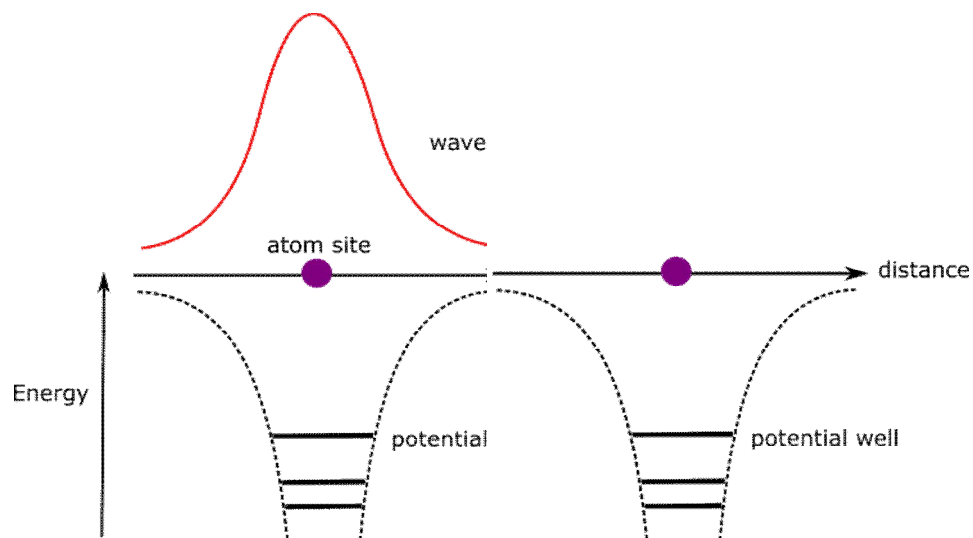


In the limit of very strong potential, Wannier functions are similar to atomic orbitals  $|\tilde{\psi}_{\vec{R},\alpha}\rangle \approx |\psi_{\vec{R},\alpha}^{atom}\rangle$

## Ansatz: Linear combination of atomic orbitals (LCOA)

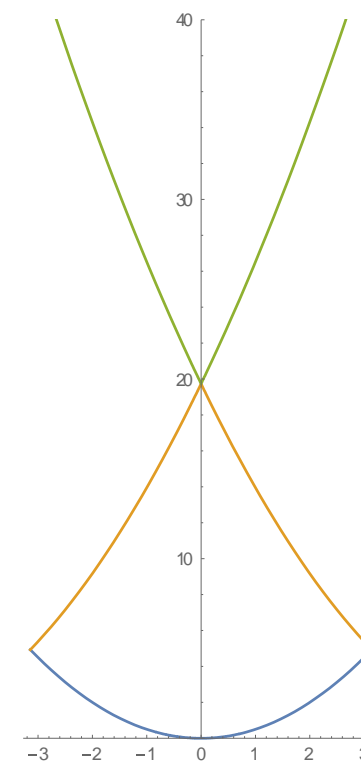
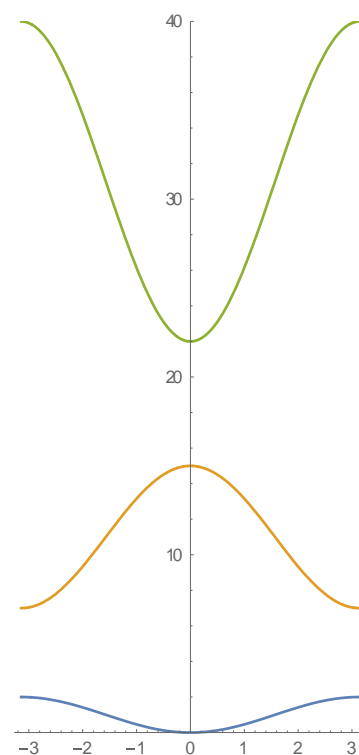
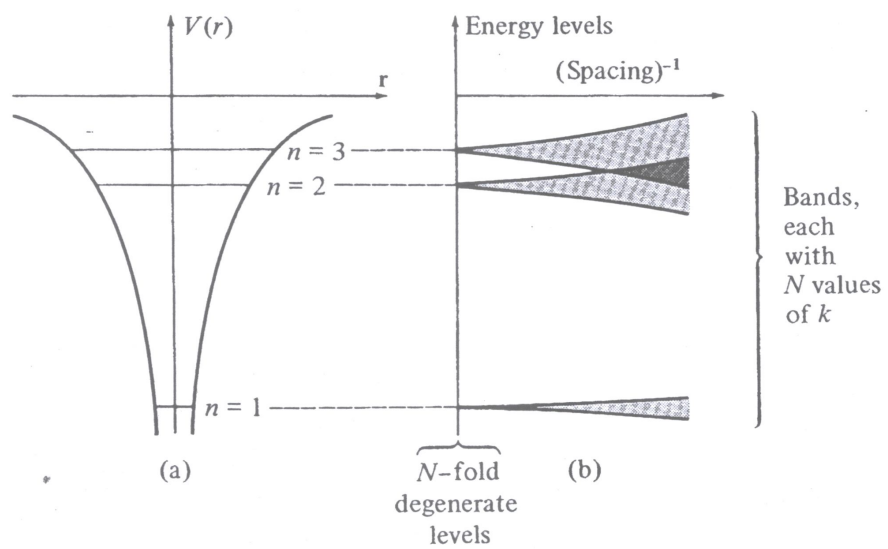
Simplest case: Hopping of one single orbital

$$t_{\alpha}(\Delta\vec{R}) \approx \langle \psi_{\vec{R}+\Delta\vec{R},\alpha}^{atom} | \hat{H} | \psi_{\vec{R},\alpha}^{atom} \rangle = \langle \psi_{\vec{R}+\Delta\vec{R},\alpha}^{atom} | (\hat{H}^{atom} + \hat{H}^{rest}) | \psi_{\vec{R},\alpha}^{atom} \rangle$$

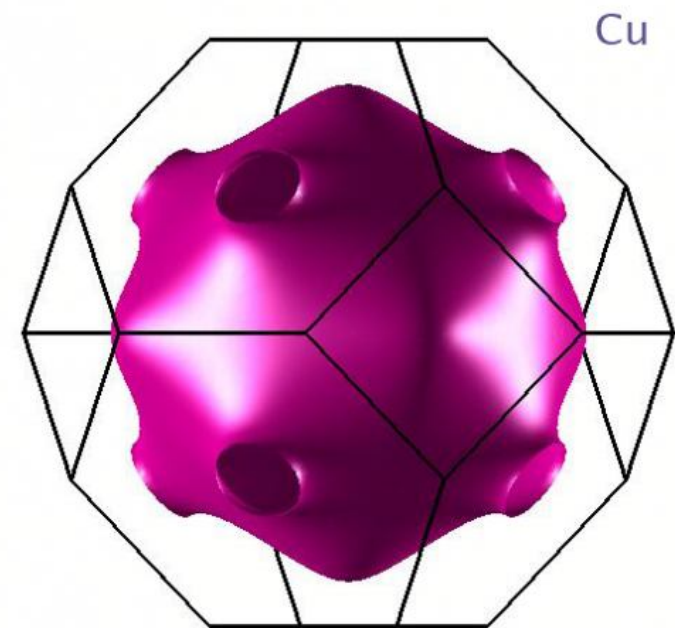
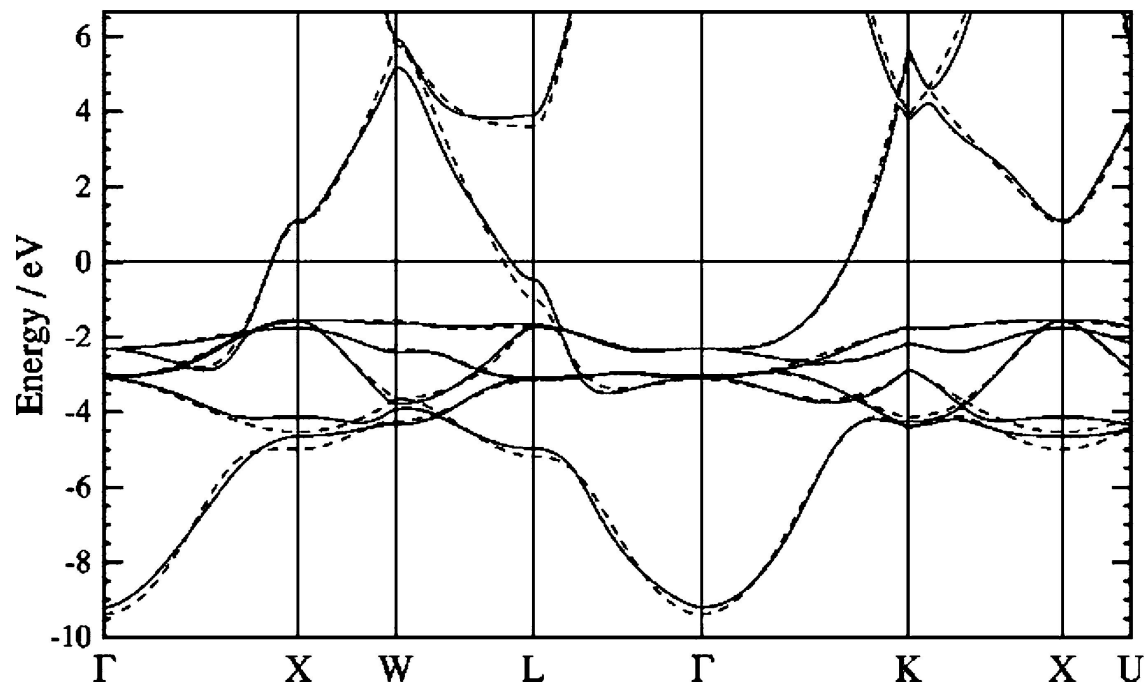


## Tight binding approximation

$$\hat{H} = \sum_{\alpha} \sum_{\vec{R}, \Delta\vec{R}} t_{\alpha}(\Delta\vec{R}) \left| \psi_{\vec{R}+\Delta\vec{R}, \alpha}^{atom} \right\rangle \left\langle \psi_{\vec{R}, \alpha}^{atom} \right|$$



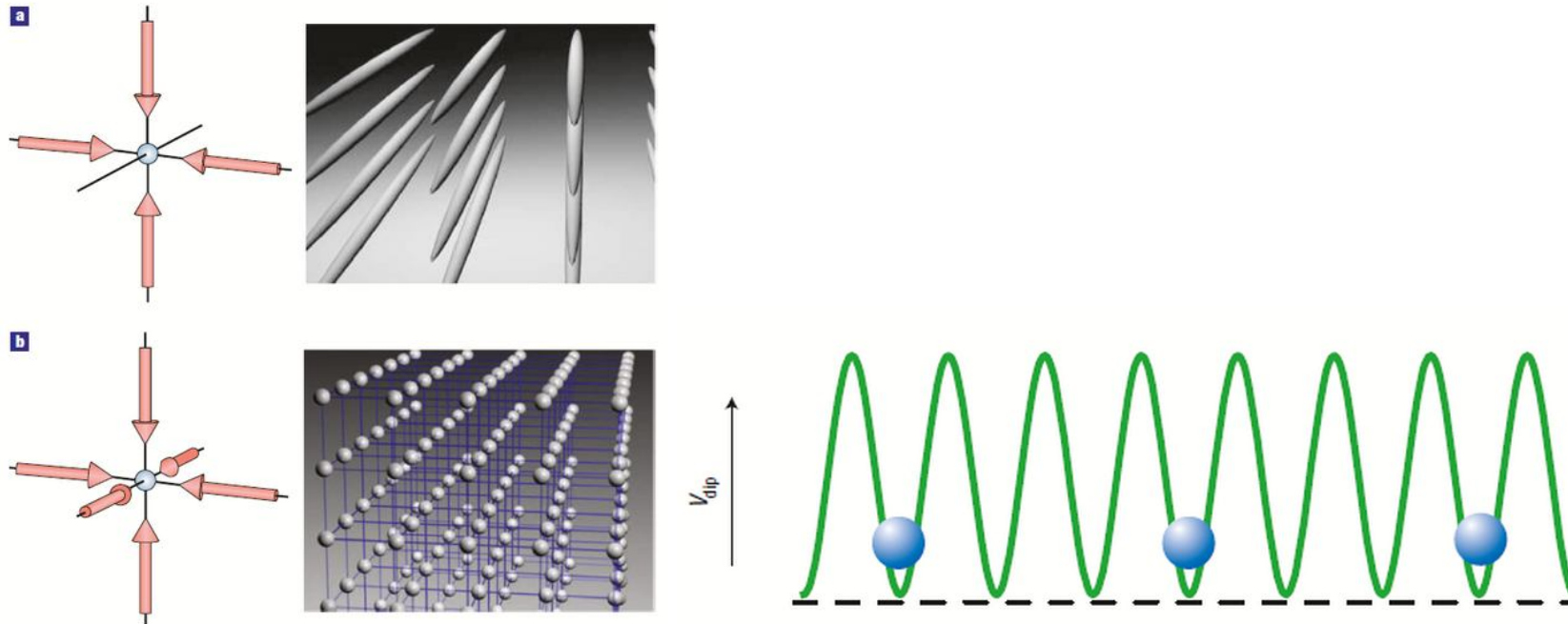
**Real band structure: Estimates**



## Real band structure in ultracold gases

Immanuel Bloch, Nature Phys. V1, p 23 (2005)

Wilhelm Zwerger, J. Opt B: Quantum Semiclass. Opt. 5, S9 (2003).



$$\begin{aligned}\hat{H} &= -\frac{\hbar^2}{2m}\nabla^2 + U(\vec{r}) \\ &= \hat{H}_x + \hat{H}_y + \hat{H}_z\end{aligned}$$

$$\hat{H}_x = -\frac{\hbar^2}{2m}\frac{\partial^2}{\partial x^2} + U \sin^2 \frac{\pi x}{a}$$

**The Mathieu equation**

$$\hat{H}_{1D} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial r^2} + U \sin^2 k_a r$$

Define  $\tau = k_a r$

Müller-Kirsten: *Introduction to Quantum Mechanics*, World Scientific  
Martin Holthaus, J. Phys. B: At. Mol. Opt. Phys. **49**, 013001 (2006)

$$\left( \frac{\partial^2}{\partial \tau^2} + a - 2p \cos \tau \right) \varphi_{1D}(\tau) = 0$$

