

If f is balanced the amplitude of $|y = 0, 0, \dots, 0\rangle$ is zero since then

$$(104) \quad \sum_{\underline{x}=0}^{2^n-1} (-1)^{(\underline{x} \cdot \underline{y} + f(\underline{x}))} = \sum_{\underline{x}=0}^{2^n-1} (-1)^{f(\underline{x})} = 0$$

If f is constant $f(\underline{x}) = f \in \{0, 1\}$

$$(105) \quad \sum_{\underline{x}=0}^{2^n-1} (-1)^{(\underline{x} \cdot \underline{y} + f(\underline{x}))} = (-1)^f \sum_{\underline{x}=0}^{2^n-1} (-1)^{\underline{x} \cdot \underline{y}}$$

$= 0$ for all $\underline{y}$ except $\underline{y} = (0, 0, 0, \dots, 0)$

□

The Deutsch-Jozsa problem is however a problem within BPP, i.e. it can be efficiently solved probabilistically with a classical computer.

3.4 BQP $\neq$ BPP : the Simon problem

We will now discuss a problem which is believed to be outside of BPP but can be solved with polynomial effort using a quantum algorithm:

Simon problem

given $f: \{0,1\}^n \rightarrow \{0,1\}^n$

with a binary period a , i.e.

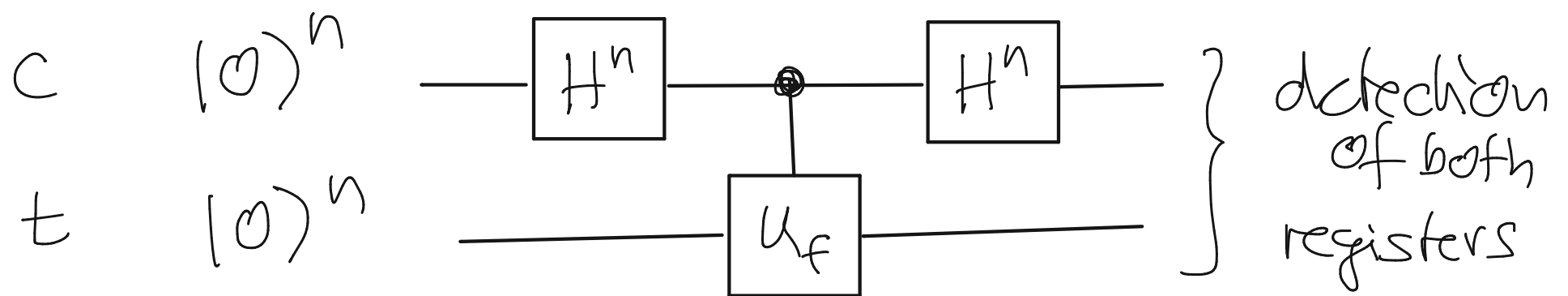
$$(106) \quad f(\underline{x}) = f(\underline{y}) \quad \text{if} \quad \underline{y} = \underline{x} \oplus \underline{a}$$

Here $\oplus$ is bitwise addition modulo 2. The task is to find the period a (note $\underline{x} \oplus \underline{a} \oplus \underline{a} = \underline{x}$)

- classically this is a NP hard problem since we have to test pairs $\underline{x}, \underline{y}$ such that

$$(107) \quad \underline{x} \oplus \underline{y} = \underline{a} \quad \text{note:} \quad (\underline{x} \oplus \underline{x} = \underline{0})$$

- quantum mechanically



where $f(\underline{x})$ is encoded in U_f

short notation $|\underline{0}\rangle \equiv |0\rangle^n = |0\rangle|0\rangle \dots |0\rangle$

$|\underline{x}\rangle, |\underline{y}\rangle$ in n -dim. computational basis
e.g. $|\underline{x}\rangle = |001101011\dots\rangle$

$$(108) \quad U_f |\underline{x}\rangle |\underline{y}\rangle = |\underline{x}, \underline{y} \oplus f(\underline{x})\rangle$$

then

$$\begin{aligned}
 & | \underline{0} \rangle | \underline{0} \rangle \xrightarrow{H^n} \frac{1}{2^{n/2}} \sum_{\underline{x}=0}^{2^n-1} | \underline{x} \rangle | \underline{0} \rangle \\
 (109) \quad & \xrightarrow{U_f} \frac{1}{2^{n/2}} \sum_{\underline{x}=0}^{2^n-1} | \underline{x} \rangle | f(\underline{x}) \rangle
 \end{aligned}$$

- measure 2nd register (target) $\Rightarrow$ this yields randomly one of the 2^n possible outputs of f :

$$\begin{array}{l}
 1 \quad | f(0,0,0,\dots) \rangle \\
 n \quad | f(1,0,0,\dots) \rangle, | f(0,1,0,\dots) \rangle \dots | f(0,0,\dots,0,1) \rangle \\
 n(n-1) \quad | f(1,1,0,\dots) \rangle, | f(1,0,1,\dots) \rangle \dots | f(0,0,\dots,1,1) \rangle \\
 \vdots \quad \quad \quad \vdots \\
 1 \quad | f(1,1,1,\dots) \rangle
 \end{array}$$

Suppose measurement yielded $| f(\underline{x}_0) \rangle$
 e.g. $| f(0,1,0,1,0,\dots,0) \rangle$

Since both $\underline{x} = \underline{x}_0$ and $\underline{x} = \underline{x}_0 \oplus \underline{a}$ give the same result, the first register is after the measurement in the state

$$(110) \quad \frac{1}{\sqrt{2}} \left(| \underline{x}_0 \rangle + | \underline{x}_0 \oplus \underline{a} \rangle \right)$$

Now how can we determine $\underline{a}$?

- bitwise Hadamard of 1st register

$$\begin{aligned}
 (111) \quad & H^n \frac{1}{\sqrt{2}} \left(|\underline{x}_0\rangle + |\underline{x}_0 \oplus \underline{a}\rangle \right) = \\
 & = \frac{1}{2^{(n+1)/2}} \sum_{\underline{y}=0}^{2^n-1} \left[(-1)^{\underline{x}_0 \cdot \underline{y}} + (-1)^{(\underline{x}_0 + \underline{a}) \cdot \underline{y}} \right] |\underline{y}\rangle \\
 & = \frac{1}{2^{(n+1)/2}} \sum_{\underline{y}=0}^{2^n-1} (-1)^{\underline{x}_0 \cdot \underline{y}} \left[1 + (-1)^{\underline{a} \cdot \underline{y}} \right] |\underline{y}\rangle
 \end{aligned}$$

Now since

$$(112) \quad \underline{a} \cdot \underline{y} = a_1 y_1 + a_2 y_2 + \dots \pmod{2} \in \{0, 1\}$$

$\Rightarrow$

$$\begin{aligned}
 (113) \quad & H^n \frac{1}{\sqrt{2}} \left(|\underline{x}_0\rangle + |\underline{x}_0 \oplus \underline{a}\rangle \right) \\
 & = \frac{1}{2^{(n+1)/2}} \sum_{\underline{a} \cdot \underline{y} = 0} (-1)^{\underline{x}_0 \cdot \underline{y}} |\underline{y}\rangle
 \end{aligned}$$

- measurement of $|\underline{y}\rangle$ register projects to one of the 2^{n-1} states $|\underline{y}\rangle$ with $\underline{a} \cdot \underline{y} = 0$

this gives one value of $\underline{y}$, say $\underline{y}_0$

- repeat the protocol until we have n linear independent binary vectors $\underline{y}_i$

with

$$(114) \quad \underline{y_j} \cdot \underline{a} = 0 \quad j=1,2,\dots,n$$

This yields n linear equations for the n components of $\underline{a}$

$$(115) \quad y_{j1}a_1 + y_{j2}a_2 + \dots + y_{jn}a_n \stackrel{\text{mod } 2}{=} 0 \quad j=1..n$$

$\Rightarrow$ determine $\underline{a}$

In order to find n linear independent vectors $\underline{y_j}$ the algorithm has to be performed on average $O(n)$ times i.e. with polynomial effort.

The Simon problem shows

$$BQP > BPP$$

(i.e. $BQP \neq BPP$)

3.5 Grover's search algorithm

Problem: given an unsorted list with N entries (reverse telephone book)
How to find an entry (e.g. a given telephone number)?

- classically $N/2$ inquiries are needed!

remark: a search in an ordered list requires only $O(\log_2 N)$ inquiries.
(successive bi-section)

$$0 \leq x \leq N$$

$$0 \leq x < N/2$$

$$N/2 \leq x \leq N$$

$$0 \leq x < N/4$$

$$N/4 \leq x < N/2$$

$$N/2 \leq x < \frac{3}{4}N$$

$$\frac{3}{4}N \leq x \leq N$$

⋮ etc

Lev Grover found an algorithm that finds an entry in

$$O(\sqrt{N})$$

steps quantum mechanically. This is practically very relevant but does not change the complexity class as

$$(116) \quad N = 2^{\log_2 N} \quad \sqrt{N} = 2^{\frac{1}{2} \log_2 N}$$

both scale exponentially with number of bits $(\log_2 N)$ needed to encode N

Let's assume that the unsorted list (inverse telephone book) is encoded in an oracle

For a given telephone number $\underline{\omega}$ corresponding to the person's name $\underline{x}_0$ we have

$$(117) \quad \begin{aligned} f_{\underline{\omega}}(\underline{x}) &= 1 & \text{for } \underline{x} = \underline{x}_0 \\ f_{\underline{\omega}}(\underline{x}) &= 0 & \text{else} \end{aligned}$$

quantity we look for!

$f_{\underline{\omega}}$ is encoded in a unitary $U_{\underline{\omega}}$

$$(118) \quad U_{\underline{\omega}} |\underline{x}\rangle |\underline{y}\rangle = |\underline{x}\rangle |\underline{y} \oplus f_{\underline{\omega}}(\underline{x})\rangle$$

Note that for given $\underline{x}$ (name) it is easy to find $\underline{\omega}_0$ (telephone number) since in this direction the list is ordered, i.e. $\underline{x} \rightarrow f_{\underline{\omega}}(\underline{x})$ is easy: if $\underline{\omega} = \underline{\omega}_0$ then $f_{\underline{\omega}} = 1$ else $f_{\underline{\omega}} = 0$.

As seen before

$$(119) \quad U_{\underline{\omega}} |\underline{x}\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) = (-1)^{f_{\underline{\omega}}(\underline{x})} |\underline{x}\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

I.e. ignoring the second register which stays unchanged, the action of $U_{\underline{\omega}}$ is

$$(120) \quad U_{\underline{\omega}} |\underline{x}\rangle \longrightarrow (-1)^{f_{\underline{\omega}}(\underline{x})} |\underline{x}\rangle$$

This means $U_{\underline{\omega}}$ changes the sign of all states orthogonal to the state $|\underline{x}_0\rangle$ corresponding to $\underline{\omega}$ (input telephone number).

(130)

$$U_w \hat{=} \mathbb{1} - 2 |\underline{x}_0\rangle \langle \underline{x}_0|$$

Now we use as input a superposition of computational basis states as input

(131)

$$|\underline{s}\rangle = \frac{1}{\sqrt{N}} \sum_{\underline{x}=0}^{N-1} |\underline{x}\rangle$$

$|\underline{s}\rangle$ contains the state $|\underline{x}_0\rangle$ which we are looking for with probability $P_{x_0} = 1/N$.

Grover's Algorithm aims at an iterative amplification of this probability. To this end we apply

(132)

$$R = \bar{U}_s U_w$$

where

(133)

$$\bar{U}_s = 2|\underline{s}\rangle \langle \underline{s}| - \mathbb{1}$$

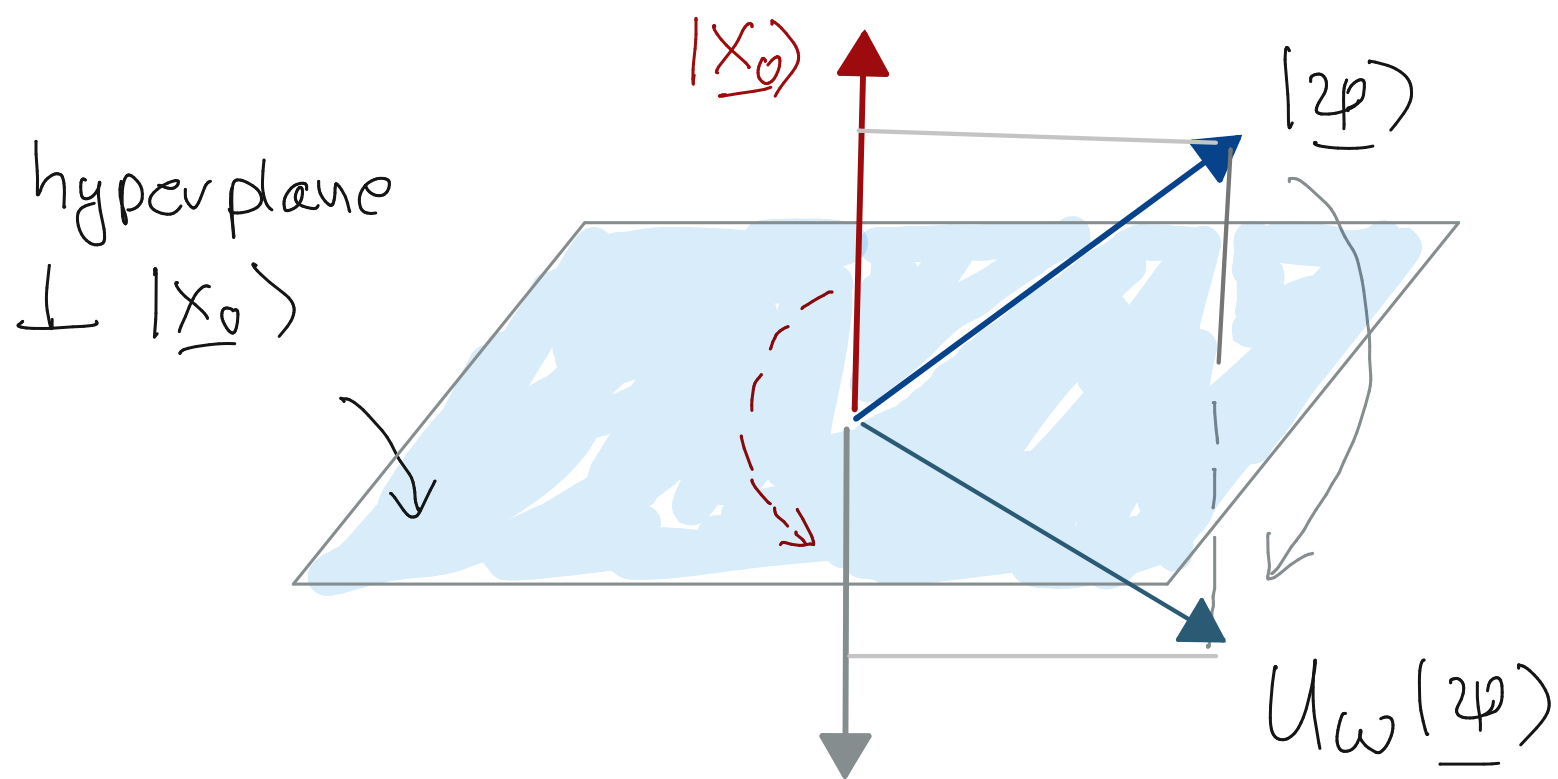
R is called reflection about the mean

geometric illustration

(134)

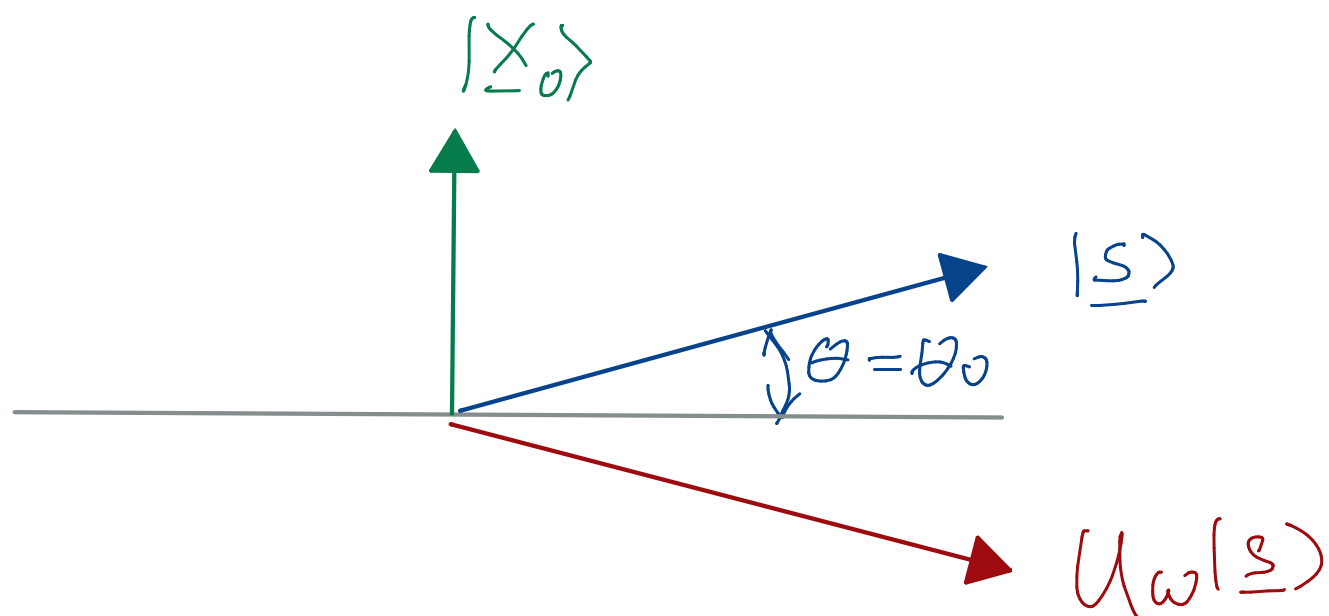
$$|\langle \underline{s} | \underline{x}_0 \rangle| = \frac{1}{\sqrt{N}} \doteq \sin \theta \approx \theta$$

- action of U_w : The component of a general state $|\varphi\rangle$ proportional to $|\underline{x}_0\rangle$ changes sign

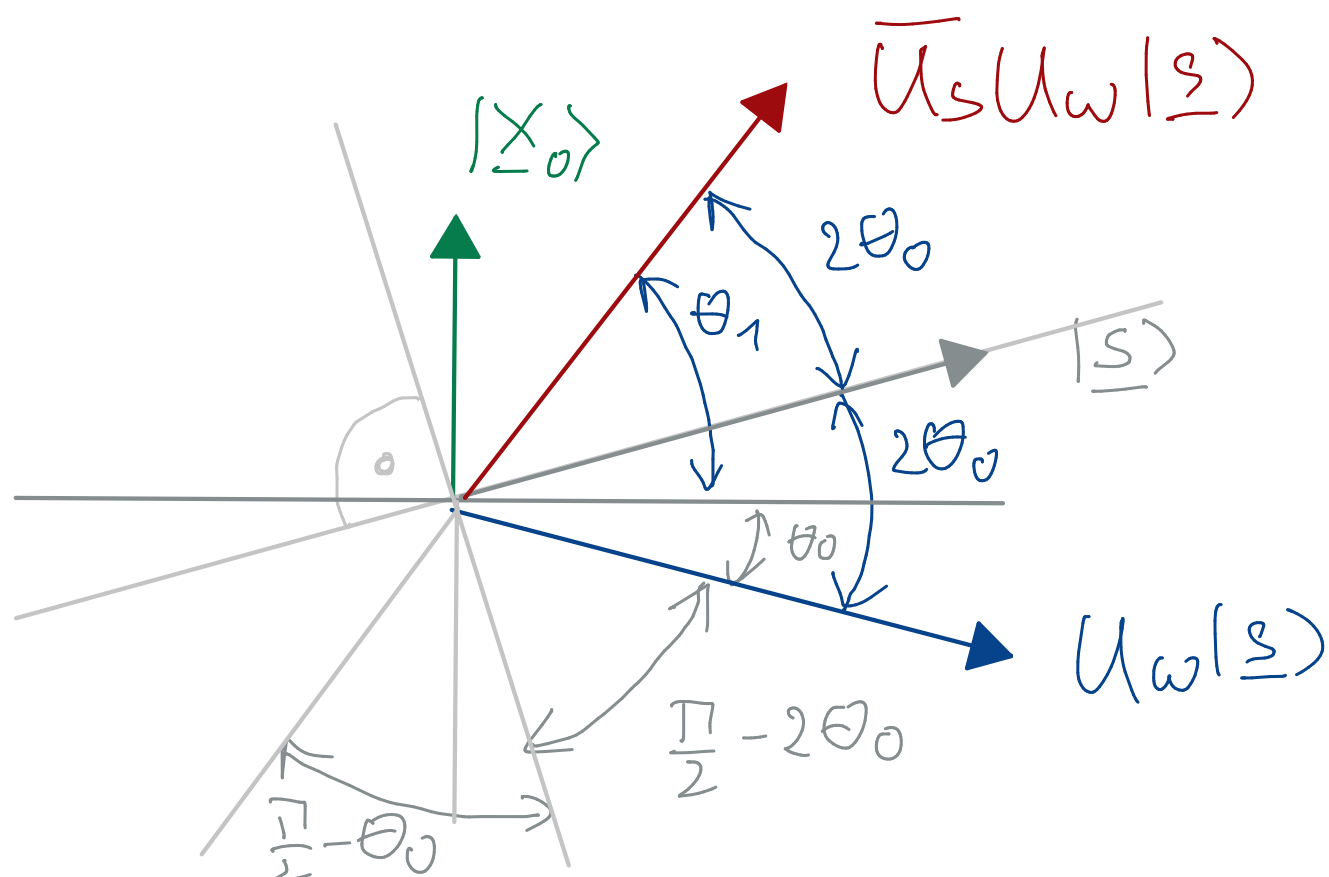


• action of $\bar{U}_S U_w$

(i)



(ii)



(135) $\theta_1 = \pi - 2\left(\frac{\pi}{2} - 2\theta_0\right) - \theta_0 = 3\theta_0$

i.e. in every iteration step the vector is shifted by $\Delta\theta = 2\theta_0$ in the direction of $|\underline{x}_0\rangle$

$$(136) \quad \text{So } |\psi_0\rangle = |\underline{x}_0\rangle \quad \text{and } |\psi_{k+1}\rangle = R |\psi_k\rangle$$

$$(137) \quad \text{Thus } |\psi_k\rangle = \alpha_k \sum_{\underline{x} \neq \underline{x}_0} |\underline{x}\rangle + \beta_k |\underline{x}_0\rangle$$

where

$$(138) \quad \alpha_k = \frac{1}{\sqrt{N-1}} \cos((2k+1)\theta) \quad \beta_k = \sin((2k+1)\theta)$$

The amplitude of the state $|\underline{x}_0\rangle$, which we are looking for keeps increasing as long as

$$(2k+1)\theta \leq \frac{\pi}{2}$$

Optimum number $k_{\text{opt}} = T$ of steps is $(2T+1)\theta = \frac{\pi}{2}$. Using that for large N $\sin\theta \simeq \theta = \frac{1}{\sqrt{N}} \Rightarrow$

$$(139) \quad \boxed{T = \frac{\pi}{4} \sqrt{N} + \mathcal{O}(1)}$$

number
of steps
in Grover
algorithm

$$\text{and } \sin((2T+1)\theta) \simeq 1.$$

I.e. after $O(\sqrt{N})$ steps $|\underline{x}_0\rangle$ is approximated with probability close to 1 !

Implementation of $\bar{U}_S$

$$(140) \quad |\underline{s}\rangle = H^{(n)} |\underline{0}\rangle = \frac{1}{2^{n/2}} \sum_{\underline{x}=0}^{2^n-1} |\underline{x}\rangle$$

i.e.

$$(141) \quad \bar{U}_S = 2|\underline{s}\rangle\langle\underline{s}| - \mathbb{I} = H^{(n)} (2|\underline{0}\rangle\langle\underline{0}| - \mathbb{I}) H^{(n)}$$

$H^{(n)}$ is easy, but how about $2|\underline{0}\rangle\langle\underline{0}| - \mathbb{I}$,
or equivalently $1 - 2|\underline{0}\rangle\langle\underline{0}|$?

$$1 - 2|\underline{0}\rangle\langle\underline{0}| = \underbrace{\bar{X}^{(n)}}_{\text{bitwise Pauli X}} \underbrace{(1 - 2|\underline{1}\rangle\langle\underline{1}|)}_{\text{changes sign of } (1111\dots 1) \text{ only}} \bar{X}^{(n)}$$

$1 - 2|\underline{1}\rangle\langle\underline{1}|$ is unitary as

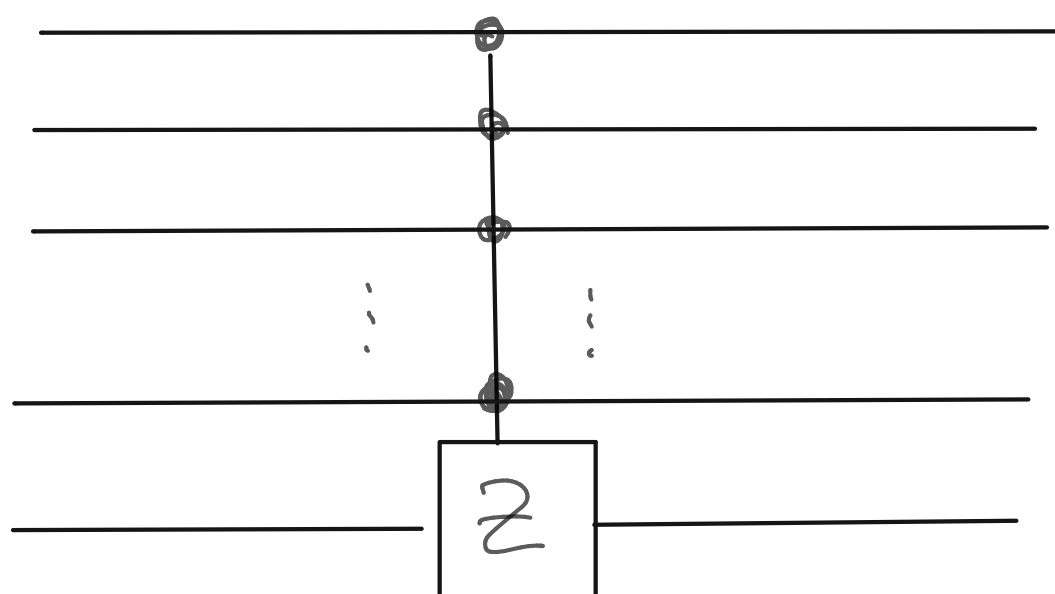
$$(1 - 2|\underline{1}\rangle\langle\underline{1}|)^{\dagger} = (1 - 2|\underline{1}\rangle\langle\underline{1}|)$$

$$\text{and } (1 - 2|\underline{1}\rangle\langle\underline{1}|)(1 - 2|\underline{1}\rangle\langle\underline{1}|) =$$

$$= 1 - 4|\underline{1}\rangle\langle\underline{1}| + 4|\underline{1}\rangle\langle\underline{1}| = 1 \quad \checkmark$$

and can be implemented by an N -fold controlled gate

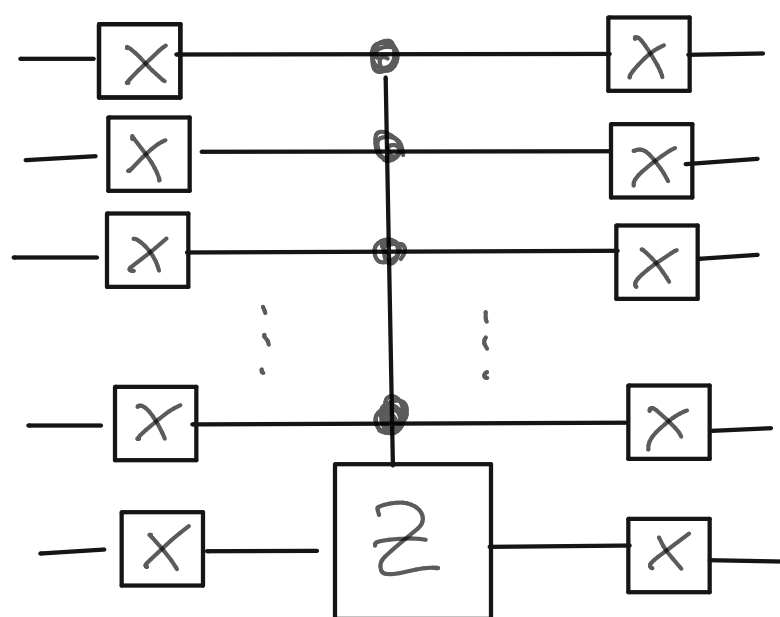
$$1 - 2(1)(1)$$



all of these bits need to be in $|1\rangle$ else gate does nothing

If all but the last qubit are in $|1\rangle$ and the last in $|0\rangle$ the state is unchanged. If the last qubit is in $|1\rangle$ the sign is flipped to, i.e.

So $\bar{U}_S$ is



This gate can be implemented with $O(n = \log N)$ elementary gates. I.e. Grover's search algorithm requires

$$(142) \quad \sim \sqrt{N} \text{ poly}(\log N)$$

gate operations.

One can show that Grover's algorithm is optimal for a search in an unsorted list