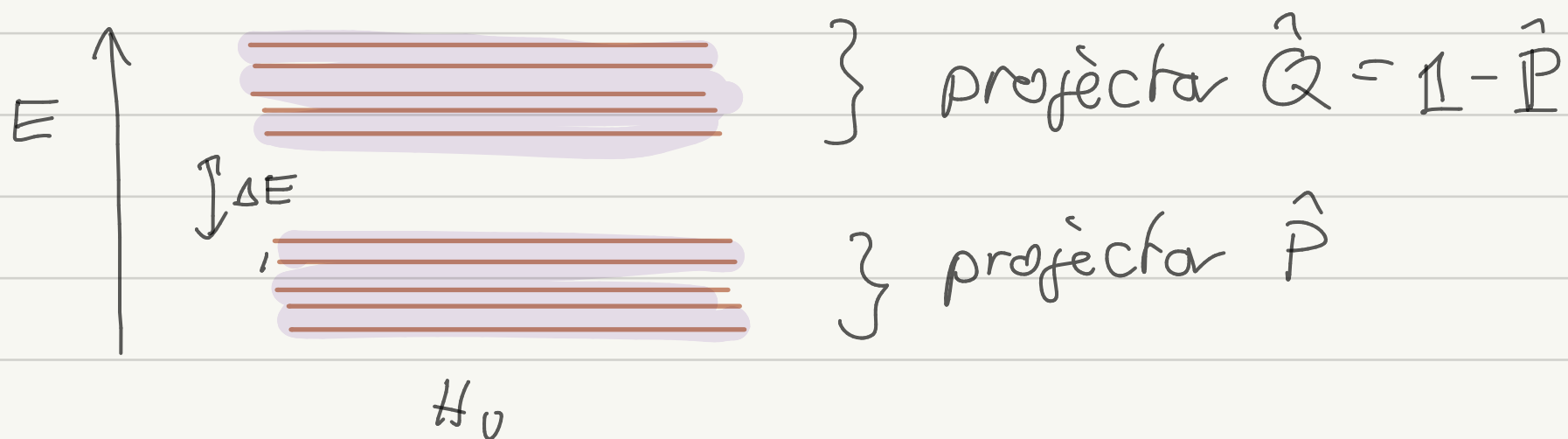


I.4 Adiabatic Elimination

Let us consider again a hamiltonian

$$(59) \quad H = H_0 + H_1$$

where the spectrum of H_0 separates in clearly distinct energy regions and H_1 is a perturbation



Consider projectors onto subspaces of H_0

$$\hat{P} \text{ and } \hat{Q} \quad \hat{P}^2 = \hat{P} \quad \hat{Q}^2 = \hat{Q}$$

Assume further that

$$(50) \quad |\hat{P} H \hat{Q}| = |\hat{P} H_1 \hat{Q}|$$

is small compared to

$$\min(|QHQ| - |PHP|) \sim \Delta E$$

If the initial state of a quantum evolution is a low energy state such that

$$(61) \quad | \psi_0 \rangle = \hat{P} | \psi_0 \rangle$$

then the time evolution remains approximately in the subspace of low-energy states spanned by $\hat{P}$ and we can derive an effective Hamiltonian

$$\begin{aligned} (62) \quad i\hbar \frac{d}{dt} \hat{P} | \psi(t) \rangle &= \hat{P} H | \psi(t) \rangle \\ &= (\hat{P} H \hat{P} + \hat{P} H \hat{Q}) | \psi(t) \rangle \\ &= P H P \underline{P | \psi(t) \rangle} + P H Q \underline{Q | \psi(t) \rangle} \end{aligned}$$

and

$$\begin{aligned} (63) \quad i\hbar \frac{d}{dt} \hat{Q} | \psi(t) \rangle &= \hat{Q} H | \psi(t) \rangle \\ &= (Q H P + Q H Q) | \psi(t) \rangle \\ &= Q H P \underline{P | \psi(t) \rangle} + Q H Q \underline{Q | \psi(t) \rangle} \end{aligned}$$

since $Q H Q$ is a high energy, we can approximate,

$$(64) \quad | Q H Q Q | \psi(t) \rangle \gg | i\hbar \frac{d}{dt} Q | \psi \rangle$$

$\Rightarrow$

$$Q | \psi(t) \rangle \simeq - (Q H Q)^{-1} Q H P P | \psi(t) \rangle$$

substituting into (62) yields

$$(65) \quad i\hbar \frac{d}{dt} P|\psi(t)\rangle = \underbrace{\left[PHP - PHQ (QHQ)^{-1} QHP \right]}_{H_{\text{eff}}} P|\psi\rangle$$

This gives an effective Hamiltonian for the slow low-energy dynamics

$$(66) \quad H_{\text{eff}} = PHP - PH_1 Q \frac{1}{QH_1 Q} QH_1 P$$

effective Hamiltonian in adiabatic approx.

where we used that P and Q are orthogonal projectors to subspaces spanned by H_0

II. Photons

Until now we have considered charged quantum particles in external, i.e. prescribed electromagnetic fields. This is not completely satisfactory

consider e.g. Lorentz force

classical physics

$$\vec{E} \equiv -\frac{\vec{F}}{q}$$

↓
quantum physics

$$? \equiv -\frac{\vec{F}}{q} \quad \text{operator?!}$$