

# Superconducting Qubits

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# Introduction

To design a qubit, or to physically realize a quantum computing system, it is important to fulfill certain requirements such as:

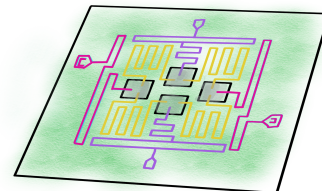
- DiVincenzo criteria (Long coherence time, scalability)
- Ease of design and manufacturing
- Ease of control

Superconducting qubits satisfy these criteria to a level

# Introduction

## Superconducting quantum computing

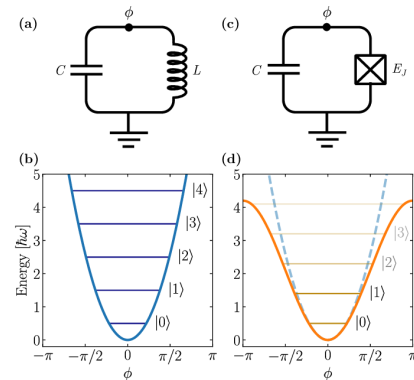
- Circuits which are designed to act as 2-level systems at low temperatures
- At this condition, the circuits have almost infinite conductivity and zero resistance.
- The charge carriers are Cooper pairs of electrons, which are bosons.
- At low temperatures, the Cooper pairs collapse to the ground state to form BE condensates
- Since there are less interactions among bosons, they act as a superconductor.



# Introduction

## Superconducting quantum computing

- The circuits are composed of semiconductor circuit elements (LC circuit, Capacitor, Inductor)
- From an LC circuit, we get the standard Quantum Harmonic oscillator with equally spaced discrete levels.
- In order to reduce to a 2-level quantum system, the Josephson junction helps with the process

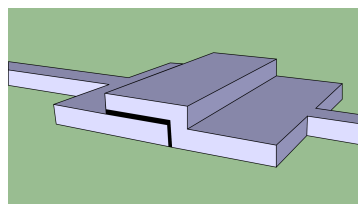
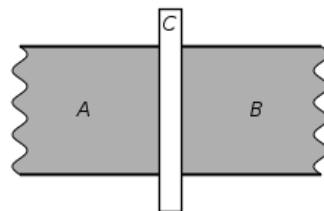


# Josephson Effect

- When 2 superconductors are restricted with a barrier, a supercurrent flows through the junction
- Given that the 2 superconductors A & B have wave functions of Cooper pairs given by

$$\psi_A = \sqrt{n_A} e^{i\phi_A}$$

$$\psi_B = \sqrt{n_B} e^{i\phi_B}$$



# Josephson Effect

The Schrodinger equation is given by

$$i\hbar \frac{\partial}{\partial t} \begin{pmatrix} \sqrt{n_A} e^{i\phi_A} \\ \sqrt{n_B} e^{i\phi_B} \end{pmatrix} = \begin{pmatrix} eV & K \\ K & -eV \end{pmatrix} \begin{pmatrix} \sqrt{n_A} e^{i\phi_A} \\ \sqrt{n_B} e^{i\phi_B} \end{pmatrix},$$

Where K is the characteristic of the junction.

Solving the equation by taking the time derivative for A,

$$\frac{\partial}{\partial t} (\sqrt{n_A} e^{i\phi_A}) = \dot{\sqrt{n_A}} e^{i\phi_A} + \sqrt{n_A} (i\dot{\phi}_A e^{i\phi_A}) = (\dot{\sqrt{n_A}} + i\sqrt{n_A} \dot{\phi}_A) e^{i\phi_A},$$

This allows to rewrite the equation as

$$\dot{\sqrt{n_A}} + i\sqrt{n_A} \dot{\phi}_A = \frac{1}{i\hbar} (eV \sqrt{n_A} + K \sqrt{n_B} e^{i\varphi}),$$

# Josephson Effect

Taking the complex conjugate and subtracting we can get the time derivative of phase  $\Phi_A'$

$$2i\sqrt{n_A} \dot{\phi}_A = \frac{1}{i\hbar} (2eV \sqrt{n_A} + K \sqrt{n_B} e^{i\varphi} + K \sqrt{n_B} e^{-i\varphi}),$$

$$\dot{\phi}_A = -\frac{1}{\hbar} (eV + K \sqrt{\frac{n_B}{n_A}} \cos \varphi).$$

Likewise, for B

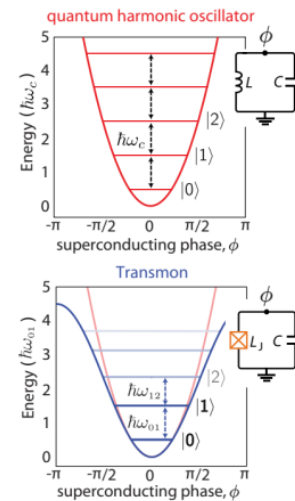
$$\dot{\phi}_B = \frac{1}{\hbar} (eV - K \sqrt{\frac{n_B}{n_A}} \cos \varphi).$$

The Josephson phase is the difference between  $\Phi_B$  and  $\Phi_A$

$$\varphi = \phi_B - \phi_A$$

# Josephson Effect

- Here, the Josephson junction with different phases on both sides causes current flow to occur by quantum tunneling
- This causes non-linear inductance which produces anharmonic oscillators with non-uniform spacing between energy levels
- Compared to QHO, this allows to access only 2 levels



# Josephson Effect

For the given phase difference, the Josephson relations are

$$I_J = I_0 \sin \phi$$

$$V = \frac{\Phi_0}{2\pi} \frac{d\phi}{dt}$$

where

$I_J$  is the Josephson current and

$\Phi_0$  is the flux quantum

# Types

Based on different degrees of freedom, there are mainly 2 categories of superconducting qubits

- Charge
- Flux

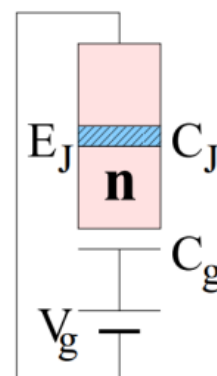
Other variants can be hybridized versions of the above 2 main types

The differentiating factor is the ratio  $E_J/EC$  where  $E_J$  is the Josephson energy and  $EC$  is the charging energy

Another means of comparison is via the number of Cooper pairs that are allowed to tunnel through the junction, therefore Charge < Flux

# Charge Qubit

- $E_J \ll EC$
- Control gate voltage  $V_g$  is coupled to the system via a gate capacitor  $C_g$
- Junction capacitance can be designed to Femtofarad range values ( $10^{-15}$  F)
- Charging energy  $E_C = e^2/2(C_g + C_J)$
- The 2 states of 0 and 1 differ by one Cooper pair charge
- Gates can be applied by tweaking the gate voltage or magnetic fields



# Charge Qubit

The Hamiltonian is given by

$$H = 4E_c(n - n_g)^2 - E_J \cos \Phi$$

$n$  = number of (excess) Cooper pairs on the island that tunnel through the junction

$\Phi$  = Josephson phase difference parameter

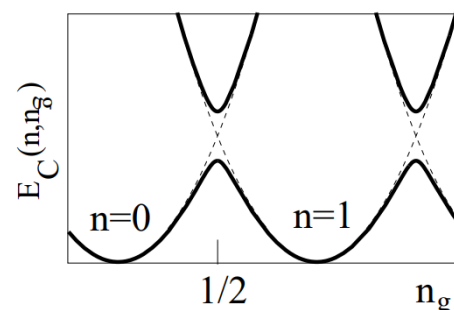
$n_g$  = gate charge which acts as an opposing control parameter =  $C_g V_g / 2e$

Considering a 2 level system and the number of Cooper pairs, the Hamiltonian rephrases to

$$\mathcal{H} = \sum_n \left\{ 4E_C(n - n_g)^2 |n\rangle \langle n| - \frac{1}{2} E_J (|n\rangle \langle n+1| + |n+1\rangle \langle n|) \right\}.$$

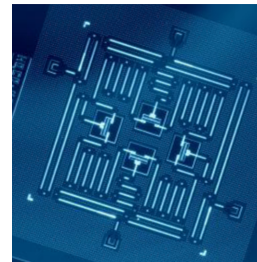
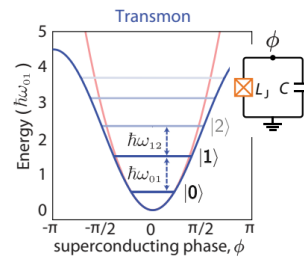
# Charge Qubit

- Initially, the system is similar to a quantum harmonic oscillator
- However, with the Josephson junctions, there is over lapping at  $\frac{1}{2}$  integer values of  $n_g$
- For the sake of 2-level system, only 2 states can be considered.
- In essence, the Josephson junctions act as valves for Cooper pairs to restrict movement within only 2 levels which is needed for a qubit



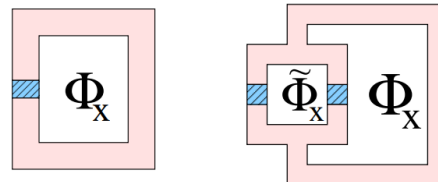
# Transmon

- As mentioned, the addition of Josephson junctions produces anharmonicity  $\alpha \equiv \omega_{12}/2\pi - \omega_{01}/2\pi$
- However, this qubit is still vulnerable to environment charge noise and other external factors
- The addition of a shunt capacitor gives us the transmon qubit
- This qubit aside from decreased charge noise sensitivity also gives increase in  $E_J/EC$  ratio and increased coherence time



# Flux Qubit

- $1 \ll E_J/EC \ll 100$
- $E_J \gg EC$
- Loops of superconducting wire fitted with a number of Josephson junctions
- Since the Josephson energy is greater, Cooper pairs flow continuously around the loop
- High coherence time ranging between 10  $\mu$ s to 23  $\mu$ s
- The only drawback is the lack of device-to device reproducibility





# Flux Qubit

The Hamiltonian is given by

$$\mathcal{H} = -E_J \cos\left(2\pi \frac{\Phi}{\Phi_0}\right) + \frac{(\Phi - \Phi_x)^2}{2L} + \frac{Q^2}{2C_J}$$

Where

$L$  is the self-inductance of the loop

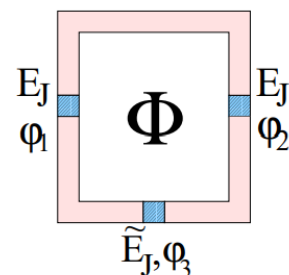
$C_J$  is the capacitance of the junction

$\Phi_x$  is the external applied flux

$Q$  is the charge on the leads which is canonically conjugate to  $\Phi$

# Flux Qubit

- When the external flux  $\Phi_x$  applied is close to a  $\frac{1}{2}$  integer number of flux  $\Phi$ , the first 2 terms of the Hamiltonian forms a double well potential where  $\Phi = \Phi_0/2$
- Since this qubit functions at low temperatures, this allows for the lowest energy states to contribute.
- The superposition of the 2 classical states from the clockwise and counterclockwise currents form the 2 level system
- Due to the fact that the product of critical current and self inductance was high, this would lead to requiring larger loop area. This could be rectified by multiple junctions

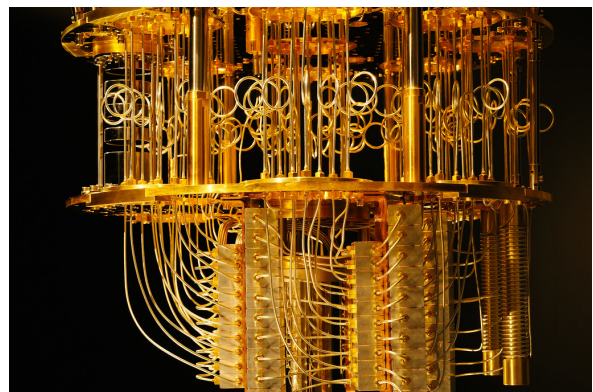


## Challenges faced

- Maintenance of high coherence
- Requires low temperatures (mK) for proper functioning, thus requiring advances in cryogenic technology first
- Fabrication of large scale circuits
- Verification/ Benchmarking of circuits becomes tougher with scale
- Requires specialized hardware that is hard to find and is also highly expensive

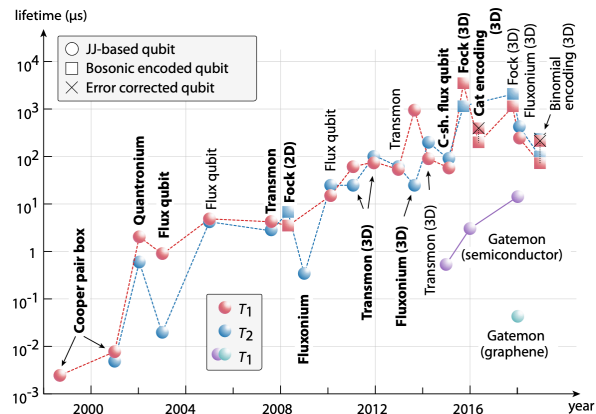
## Challenges faced

- Dilution refrigerators are custom-designed by few companies
- Helium-3 is required to operate at supercooled temperatures.
- Superconducting cables are specially designed and manufactured by only one Japanese company (Coax Co.)

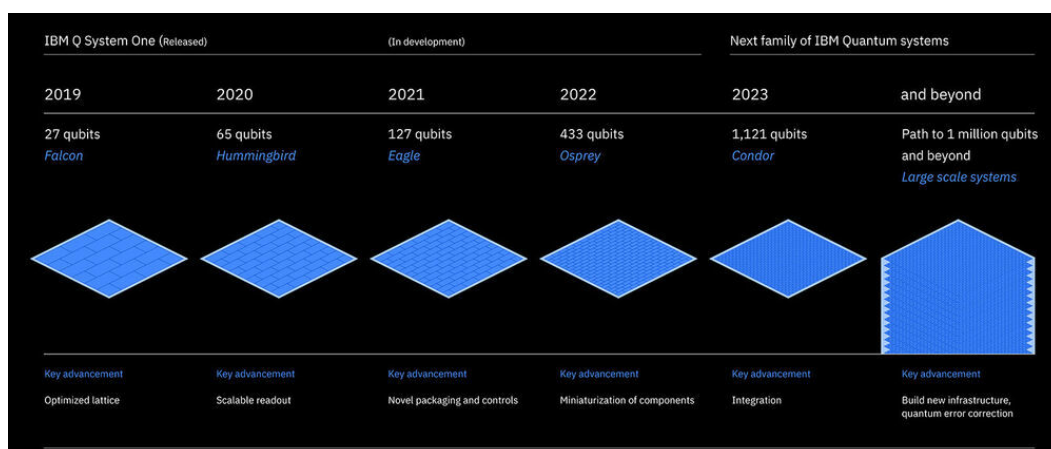


## Current progress

- Over time, the coherence and thus lifetime of various iterations of SC qubits has shown to improve with further research and development
- While there are a lot of challenges to be solved, currently SC quantum computing is seen as the most common manufacturing approach taken by companies
- Industries such as Google, IBM, D-Wave are implementing such means of quantum computing



## Current progress



IBM roadmap for Quantum processor

# References

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